

Discussion: LFUV, what's next?

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$b \rightarrow sll$ **2018 Workshop**

⇒ Assume Nature violates lepton flavour universality (muons vs electrons).

Types: Ratios of BR, based on angular observables, ...

$$\langle R_X \rangle = \frac{\langle B \rightarrow X \mu^+ \mu^- \rangle}{\langle B \rightarrow X e^+ e^- \rangle} \quad \langle Q_i \rangle = \langle P_i^\mu \rangle - \langle P_i^e \rangle \quad \langle B_k \rangle = \frac{\langle J_k^\mu \rangle}{\langle J_k^e \rangle} - 1, \quad k = 5, 6s \quad \dots$$

Properties: LFUV observables fulfill the following criteria:

- All LFUV are EXCELLENT DISCOVERY probes: $R_{X=K, K^*, \phi}$, Q_i , B_i , T_i, \dots
- Sensitivity only to the short distance part of C_9 (**charm free** in the SM).
→ only QED corrections naively $\frac{\alpha}{\pi} \log^2(\frac{m_\ell^2}{q^2})$ maybe relevant but controlled [[Bordone, Isidori, Pattori](#)]
- Sensitivity to Wilson coefficients other than C_9 .
- In presence of New Physics can be classified according to their hadronic uncertainty sensitivity:

Reminder: $\delta \mathbf{BR} \propto \delta \xi_{\perp, \parallel} +$ sensitivity to charm (known and unknown $(\delta C_9^{c\bar{c}})^2$) + p.c.

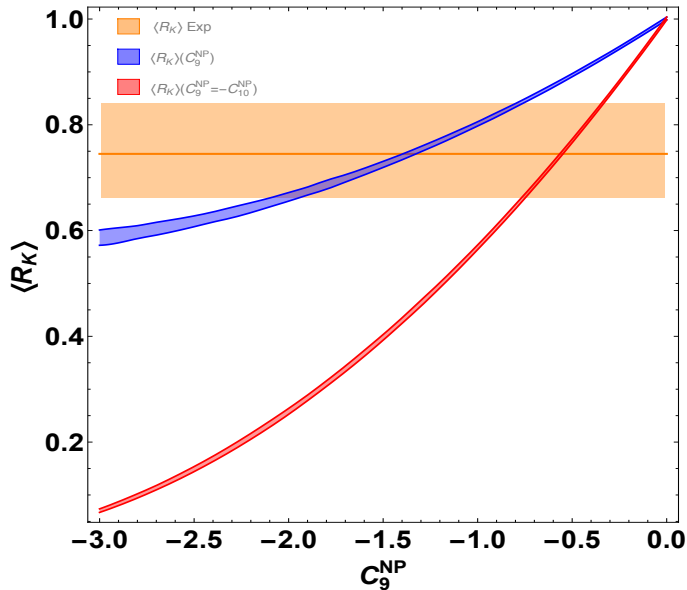
... crucial dependence on form factor error estimate ⇒ \mathbf{R}_X

$$\delta \mathbf{P}_i, \delta \mathbf{J}_{i=5,6s}^\mu / \mathbf{J}_{i=5,6s}^e \propto \alpha_s \delta \xi_{\perp, \parallel} + \text{sensitivity to charm} + \text{p.c.} \Rightarrow \mathbf{Q}_i, \mathbf{B}_i$$

Classification according to sensitivity to hadronic uncertainties in presence of New Physics: Ratios of BR, R_K

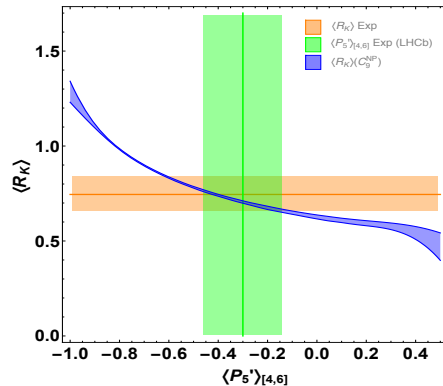
R_K : Simple structure: $f_{+,0,T} \rightarrow$ one SFF (f_+) at large-recoil.

$\rightarrow f_0$ lepton mass suppressed or arises in the presence of (pseudo)scalar while f_T suppressed by C_7^{eff} .



BLUE C_9^{NP} , RED $C_9^{\text{NP}} = -C_{10}^{\text{NP}}$

- $C_9^{\text{NP}} < 0$ and $C_{10}^{\text{NP}} > 0$ same weight adds coherently.
- Central value of R_K prefers a large negative contrib. to C_9^{NP} in excellent agreement with P_5' anomaly.

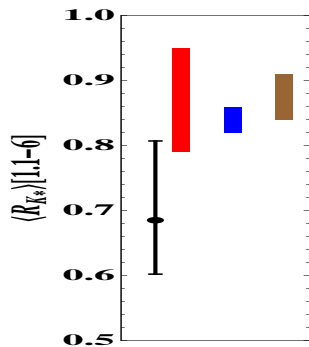


Classification according to sensitivity to hadronic uncertainties in presence of New Physics: Ratios of BR, R_{K^*}

R_{K^*} : More complex structure, 6-8 Amplitudes and 7 form factors.

Impact of long-distance charm from KMPW on $B \rightarrow K^*$ larger than on $B \rightarrow K$.

- In presence of NP or for $q^2 < 1 \text{ GeV}^2$ **hadronic uncertainties return**.
- The largest the NP \rightarrow the largest the uncertainty: $\delta R_{K^*} \propto (C_i^\mu - C_i^e)\delta FF$



Bins	Predictions R_{K^*}		
	[0.045, 1.1]	[1.1, 6.]	[15., 19.]
Standard Model	0.916 ± 0.025	1.000 ± 0.006	0.998 ± 0.001
$C_{9\mu}^{\text{NP}} = -1.1$	0.897 ± 0.044	0.867 ± 0.076	0.788 ± 0.005
$C_{9\mu}^{\text{NP}} = -1.76$	0.897 ± 0.084	0.832 ± 0.141	0.698 ± 0.009
$C_{9\mu}^{\text{NP}} = -C_{10\mu}^{\text{NP}} = -0.62$	0.866 ± 0.057	0.751 ± 0.027	0.714 ± 0.006

- 1st bin is expected to be SM-like.
- $C_9 < 0$ gets near saturation at large-recoil and $C_9 < 0$ $C_{10} > 0$ adds coherently.

At the point $C_{9\mu}^{\text{NP}} = -1.1$, $C_{9e}^{\text{NP}} = 0$:

KMPW-sch.1:

$$\xi_{\perp} = 0.31_{-0.10}^{+0.20}, \xi_{\parallel} = 0.10_{-0.02}^{+0.03}$$

BSZ-sch.1

$$\xi_{\perp} = 0.32 \pm 0.03, \xi_{\parallel} = 0.12 \pm 0.02$$

JC-sch.2

$$\xi_{\perp} = 0.31 \pm 0.04, \xi_{\parallel} = 0.10 \pm 0.02$$

Classification according to sensitivity to hadronic uncertainties in presence of New Physics: Different type of observables, Q_i, B_i, T_i

[CDMV'16]

Observables based on angular observables with cancelation of SFF at LO in presence of NP.

I) **Difference of Optimized observables:** $Q_i = P_i^\mu - P_i^e$.

→ *LFUV properties+ Inheritate the excellent properties of optimized observables in presence of NP*

II) **Ratios of coefficients** of angular distribution: $B_i = J_i^\mu / J_i^e - 1$ with $i=5,6s$.

→ *Linear sensitivity to δC_9 but kinematically suppressed*

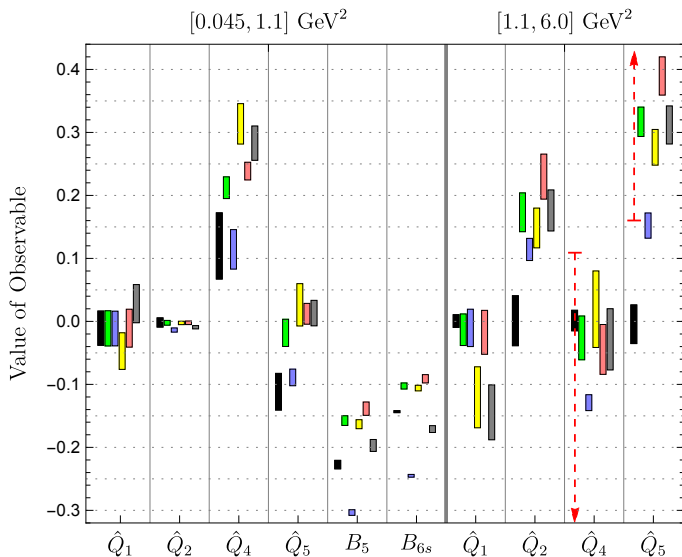
→ *Sensitivity to δC_{10} at very low- q^2*

III) **Ratios of non-optimized observables:** $T_i = \frac{S_i^\mu - S_i^e}{S_i^\mu + S_i^e}$

→ *Reduced but not perfect cancellation*

Q_i observables: Differences of Optimized observables (long bin)

[1-GREEN]: $C_{9\mu}^{\text{NP}} = -1.1$, [2-BLUE]: $C_{9\mu}^{\text{NP}} = -C_{10\mu}^{\text{NP}} = -0.61$, [3-YELLOW]: $C_{9\mu}^{\text{NP}} = -C'_{9\mu} = -1.01$,
 [4-ORANGE]: $C_{9\mu}^{\text{NP}} = -3C_{9e}^{\text{NP}} = -1.06$, [5-GRAY]: The b.f.p in the 6D. SM-[BLACK] and dashed-red [BELLE]



$$Q_i = P_i^\mu - P_i^e$$

1st bin of Q_i (not much sensitivity):

- ▶ $Q_{1,2}$ control observables.
- ▶ Q_4 only exception of sc.3 (no R_K) and sc. 6

2nd bin of Q_i :

- ▶ Q_2 useful to signal LFUV,
...not discriminating power.
- ▶ Q_5 largest discriminating power: sc. 1 & 2.
 C_{10}^{NP} subleading, sc. 1 & 2 reversed w.r.t. R_{K^*} .

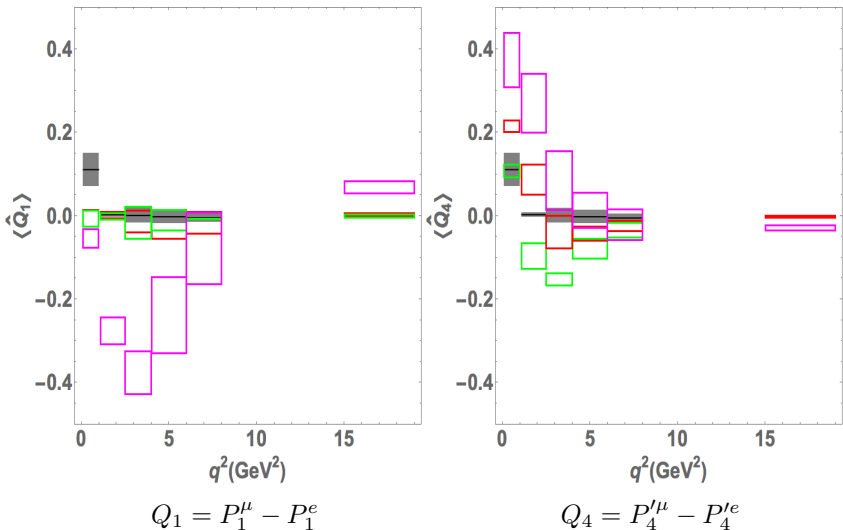
1st bin of $B_{5,6s}$ promising to disentangle C_9 from $C_9 = -C_{10}$

Probing right-handed currents (RHC) with Q_i (short bins): Different information of short versus long bins

SM predictions (grey boxes),

NP: $C_{9,\mu}^{\text{NP}} = -1.11$ & $C_{9,\mu}^{\text{NP}} = -C_{10,\mu}^{\text{NP}} = -0.65$ & $C_{9,\mu}^{\text{NP}} = -C_{9,\mu}^{\prime\text{NP}} = -1.18$ & $C_{10,\mu}^{\text{NP}} = C_{10,\mu}^{\prime\text{NP}} = 0.38$.

with $\delta C_i = C_{i,\mu} - C_{i,e}$ (and $C_{i,e}$ SM)



⇒ $Q_{1,4}$ provide excellent opportunities to probe RHC in $C'_{9,\mu}$ & $C'_{10,\mu}$.

■ Q_1 shows significant deviations in presence of RHC. If $C'_7 = 0$ at LO

$$s_0^{LO} = -2 \frac{C_7 \delta C'_{9,\mu} m_b M_B}{C_{10,\mu} \delta C'_{10,\mu} + \text{Re} C_{9,\mu} \delta C'_{9,\mu}}$$

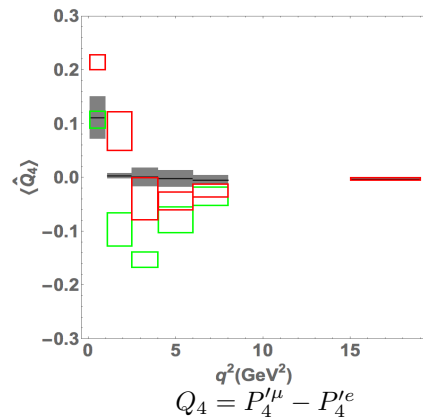
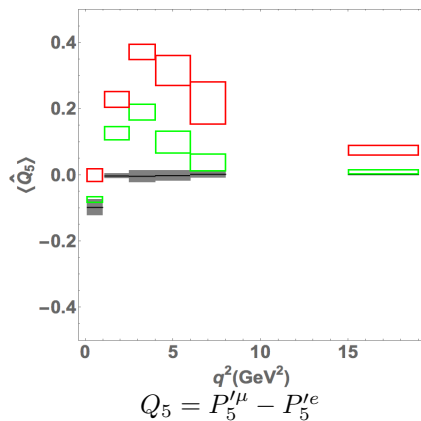
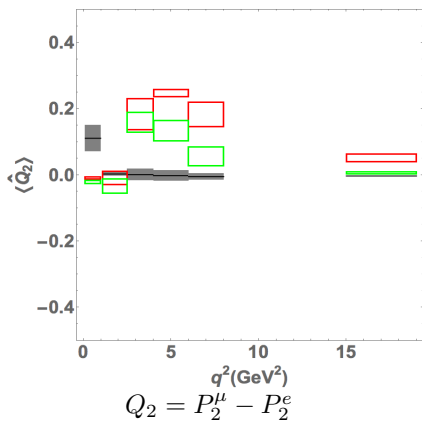
no zero (except $s = 0$) if $\delta C'_{9,\mu} = 0$.
no sensitivity to C_i if $C'_i = 0$.

■ Q_4 at low- q^2 exhibits deviations for $C'_{9,10,\mu}$ when accurate precision in measurements is achieved.

Probing NP in $C_{9,10}$ with Q_i (short bins)

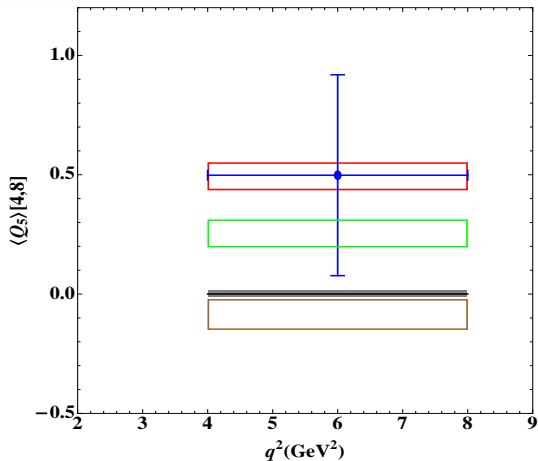
SM predictions (grey boxes),

NP: $C_{9,\mu}^{\text{NP}} = -1.11$ (scenario 1) & $C_{9,\mu}^{\text{NP}} = -C_{10,\mu}^{\text{NP}} = -0.65$ (scenario 2) with $\delta C_i = C_{i,\mu} - C_{i,e}$ (and $C_{i,e}$ SM)



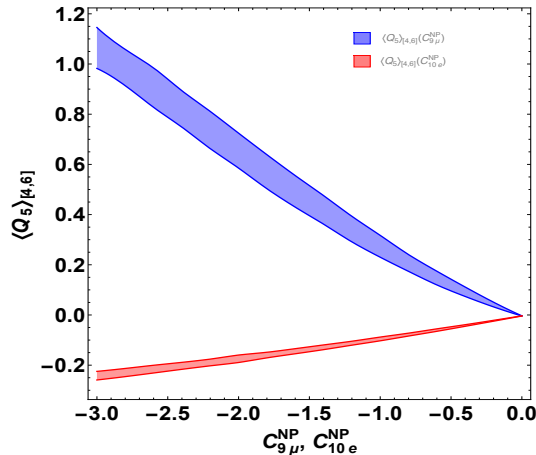
⇒ Q_2 , Q_4 & Q_5 show distinctive signatures for the two NP scenarios considered.

- Differences in the high- q^2 bins of the large recoil region of Q_2 & Q_5 are quite significant. Lack of difference between scenario 2 and SM same reason why P_5^{\prime} in scenario 2 is worst than scenario 1.
- Q_4 at very low- q^2 (second bin) is very promising to disentangle scenario 1 from 2.



BLUE Belle data.

- Green $C_{9\mu}^{\text{NP}} = -1.1$, $C_{9e}^{\text{NP}} = 0$ predicts Q_5 positive and very large.
- Red $C_{9\mu}^{\text{NP}} = -1.76$, $C_{9e}^{\text{NP}} = 0$ predicts Q_5 positive and large.
- Brown $C_{10e} < 0$, predicts Q_5 negative.



- Blue (option 1)
 $C_{9\mu}^{\text{NP}} < 0$, $C_{9e}^{\text{NP}} = 0$
predicts $Q_5 > 0$.
- Red (option 2)
 $C_{10e}^{\text{NP}} < 0$, $C_{9\mu}^{\text{NP}} = 0$
predicts $Q_5 < 0$.

If no RHC $0 \lesssim Q_5 \lesssim 0.2$ favours mechanisms with NP in $C_{9\mu} = -C_{10\mu} < 0$ or $C_{9\mu} < 0$ while $Q_5 \geq 0.2$ favours mechanisms with $C_{9\mu} < 0$, explains $P_5^{\prime\mu}$ and excludes h_λ^i explanation,

on the contrary

$Q_5 < 0$ would favour the solution $C_{10e} < 0$ and require $h_\lambda^i \neq 0$ to explain P_5^{\prime} .

Idea: Combine J_i^μ & J_i^e to build combinations sensitive to some C_i , with controlled sensitivity to long-distance charm.

$$\beta_\ell J_5^\mu - 2iJ_8 = 8\beta_\ell^2 N^2 m_B^2 (1 - \hat{s})^3 \frac{\hat{m}_{K^*}}{\hat{s}\sqrt{\hat{s}}} C_{10}^\ell \left[C_7 \hat{m}_b (1 + \hat{s}) + \hat{s} C_9^\ell \right] \xi_\perp \xi_\parallel + \dots$$

$$\beta_\ell J_{6s}^\mu - 2iJ_9 = 16\beta_\ell^2 N^2 m_B^2 \frac{(1 - \hat{s})^2}{\hat{s}} C_{10}^\ell \left[2C_7 \hat{m}_b + \hat{s} C_9^\ell \right] \xi_\perp^2 + \dots$$

Assuming real NP & maximal LFUV μ vs e , natural combinations are

$$B_5 = \frac{J_5^\mu}{J_5^e} - 1 \quad B_{6s} = \frac{J_{6s}^\mu}{J_{6s}^e} - 1$$

- Full charm insensitive in the SM.

- Linear sensitivity to δC_9 **kinematically suppressed**.

In the large-recoil limit and in absence of RHC currents [CDMV'16]:

$$B_5 = \frac{J_5^\mu - J_5^e}{J_5^e} = \frac{\beta_\mu^2 - \beta_e^2}{\beta_e^2} + \frac{\beta_\mu^2}{\beta_e^2} \frac{\delta C_{10}}{C_{10}} + \frac{\beta_\mu^2}{\beta_e^2} \frac{2(C_{10} + \delta C_{10})\delta C_9 \hat{s}}{C_{10} (2C_7 \hat{m}_b (1 + \hat{s}) + (2C_9 + \Delta C_{9,0} + \Delta C_{9,\perp}) \hat{s})} + \dots$$

$$B_{6s} = \frac{J_{6s}^\mu - J_{6s}^e}{J_{6s}^e} = \frac{\beta_\mu^2 - \beta_e^2}{\beta_e^2} + \frac{\beta_\mu^2}{\beta_e^2} \frac{\delta C_{10}}{C_{10}} + \frac{\beta_\mu^2}{\beta_e^2} \frac{2(C_{10} + \delta C_{10})\delta C_9 \hat{s}}{C_{10} (4C_7 \hat{m}_b + (2C_9 + \Delta C_{9,\perp} + \Delta C_{9,\parallel}) \hat{s})} + \dots$$

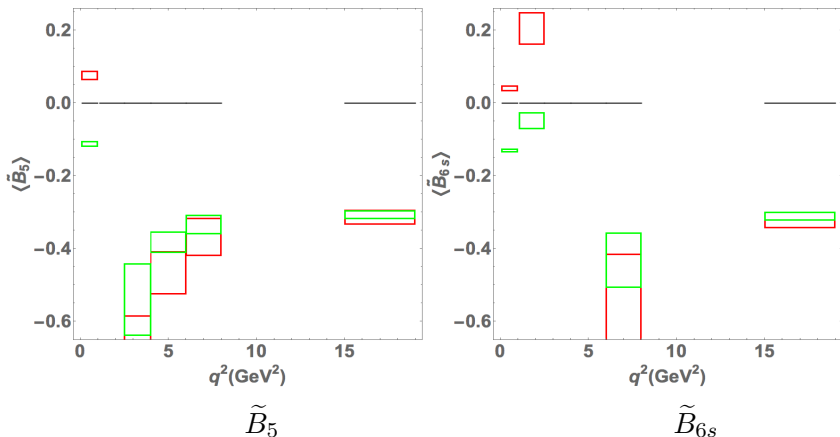
B_5 & B_{6s} are **not identically 0** in the SM.

Lepton mass differences generates a non-zero contribution mainly in the first bin.

\Rightarrow If on an event-by-event basis experimentalist can measure $\langle J_i^\mu / \beta_\mu^2 \rangle$:

$$\langle \widetilde{B}_5 \rangle = \frac{\langle J_5^\mu / \beta_\mu^2 \rangle}{\langle J_5^e / \beta_e^2 \rangle} - 1 \quad \langle \widetilde{B}_{6s} \rangle = \frac{\langle J_{6s}^\mu / \beta_\mu^2 \rangle}{\langle J_{6s}^e / \beta_e^2 \rangle} - 1$$

- SM Predictions: $\langle \widetilde{B}_i \rangle = 0.00 \pm 0.00$.
- All good properties of $B_{5,6s}$ + simpler structure $\beta_i \rightarrow 1$.



- When $\hat{s} \rightarrow 0$,
 $\widetilde{B}_5 = \widetilde{B}_{6s} = \delta C_{10} / C_{10} \Rightarrow$
Sensitivity to δC_{10} !
 Exactly as B_5, B_{6s} but simpler.
- 1st Bins: Capacity to distinguish $C_{9,\mu}^{\text{NP}} = -1.11$ from $C_{9,\mu}^{\text{NP}} = -C_{10,\mu}^{\text{NP}} = -0.65$.

BACK-UP

How LFUV NP enter in Wilson coefficients?:

$$C_{i,\mu} = \begin{cases} C_i + \delta C_i, & i = 10, 9', 10' \\ C_9 + \delta C_9 + \Delta C_9^{(j)} \end{cases} \quad C_{i,e} = \begin{cases} C_i, & i = 10, 9', 10' \\ C_9 + \Delta C_9^{(j)} \end{cases}$$
$$j = \perp, \parallel, 0$$

Notice $C_{7,7'}$ is obviously lepton-mass independent.

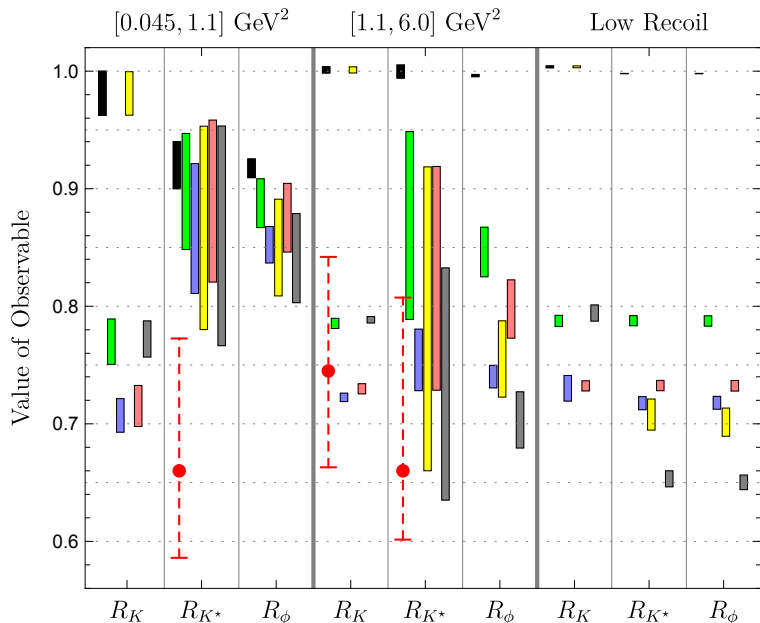
$\Rightarrow \delta C_i = C_{i,\mu} - C_{i,e} \equiv$ amount of LFU violation.

$\Rightarrow C_i \equiv$ SM + LFU NP.

$\Rightarrow \Delta C_9^{(j)} \equiv$ long-distance charm. Two types:

- **Transversity Dependent:** $\Delta C_9^{\perp, \parallel, 0}$ different.
- **Transversity Independent:** $\Delta C_9^{\perp} = \Delta C_9^{\parallel} = \Delta C_9^0$.

Disentangling New Physics: Ratios of Branching Ratios



SM-[BLACK]

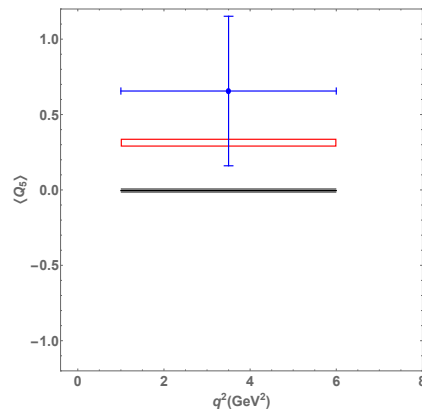
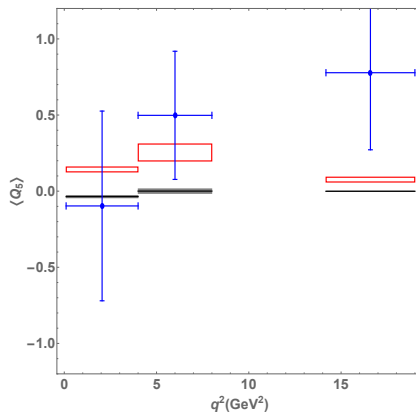
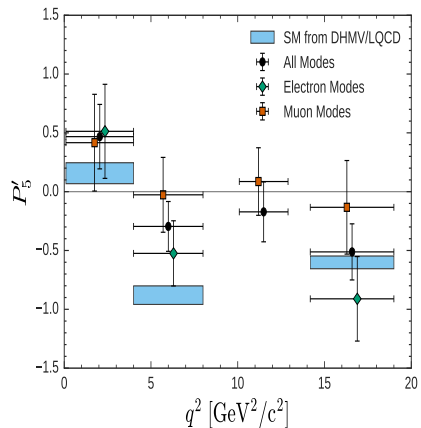
Five “good” scenarios:

- ▶ Sc. 1 [GREEN]: $C_{9\mu}^{\text{NP}} = -1.1$,
- ▶ Sc. 2 [BLUE]: $C_{9\mu}^{\text{NP}} = -C_{10\mu}^{\text{NP}} = -0.61$,
- ▶ Sc. 3 [YELLOW]: $C_{9\mu}^{\text{NP}} = -C'_{9\mu} = -1.01$,
- ▶ Sc. 4 [ORANGE]: $C_{9\mu}^{\text{NP}} = -3C_{9e}^{\text{NP}} = -1.06$,
- ▶ Sc. 5:[GRAY]: The best fit point in the six-dimensional fit.

R_{K^*} is computed using very conservative KMPW-FF but R_ϕ using BSZ-FF (only available).

ATTENTION: In presence of NP $R_{K,K^*,\phi}$ are largely sensitive to FF choices

First hint from Belle?

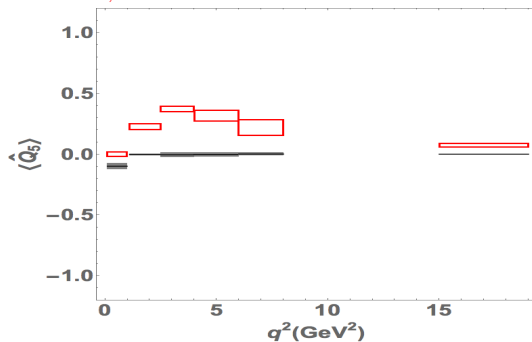
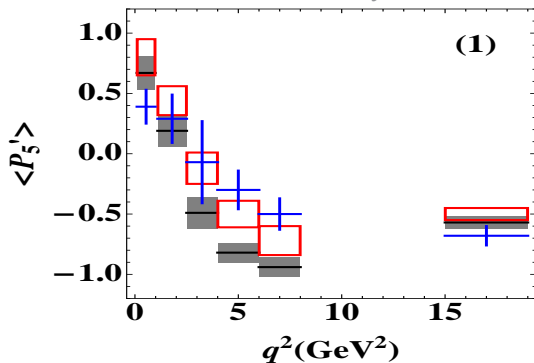


- different systematics than LHCb (combination of channels).
- Belle has found for $\langle P'_5 \rangle_{[4,8]}^\mu$ a 2.6σ deviation while 1.3σ for $\langle P'_5 \rangle_{[4,8]}^e$
- Q_5 points in the same direction as $C_{9\mu}^{\text{NP}} = -1.1$ scenario (in red).

More data needed for confirmation...

Q_i observables. The example: P'_5 versus $Q_5 = P_5^{\prime\mu} - P_5^{\prime e}$

Gray-SM, Red-NP $C_{9,\mu}^{\text{NP}} = -1.11$, $C_{9,e}^{\text{NP}} = 0$ and blue data



- Soft FF independent at LO exactly in SM
Soft FF independent at LO exactly in NP.
- Large sensitivity to $C_{9,\mu}$. SM (DHMV'15):

$$\langle P'_5 \rangle_{[4,6]} = -0.82 \pm 0.08$$

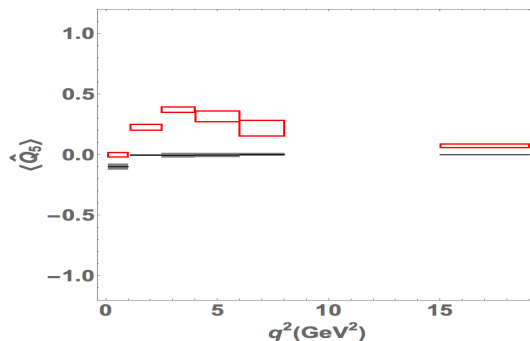
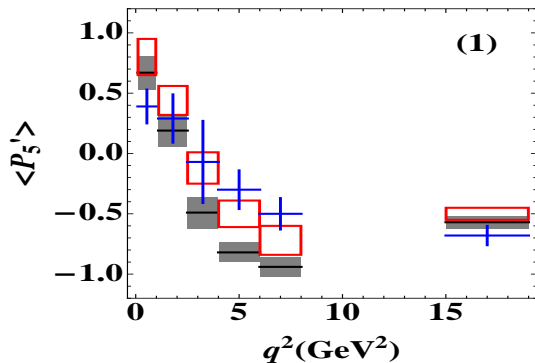
$$\langle P'_5 \rangle_{[6,8]} = -0.94 \pm 0.08$$

- FF indep. at all orders in SM (up to $\Delta m_\ell^2/q^2$).
Soft FF indep. at LO exactly in NP.
- Long-distance charm insensitive in the SM.
Large sensitivity to $\delta C_9 = C_{9,\mu} - C_{9,e}$.
(CDMV'16): ($< 10^{-3}$ without lepton mass)

$$\langle \hat{Q}_5 \rangle_{[4,6]} = -0.002 \pm 0.017$$

$$\langle \hat{Q}_5 \rangle_{[6,8]} = +0.002 \pm 0.010$$

Q_i observables. The example: P'_5 versus $Q_5 = P'_5{}^\mu - P'_5{}^e$ for $C_{9,\mu}^{\text{NP}} = -1.1$



Remark: In presence of NP hadronic uncertainties reemerge in \hat{Q}_5 ...

P'_5	Prediction $C_{9,\mu}^{\text{NP}} = -1.1$	\hat{Q}_5	SM-Prediction	\hat{Q}_5	Prediction $\delta C_9 = -1.1$
[0.1, 0.98]	$+0.80 \pm 0.14$	[0.1, 0.98]	-0.097 ± 0.023	[0.1, 0.98]	0.000 ± 0.018
[1.1, 2.5]	$+0.43 \pm 0.12$	[1.1, 2.5]	-0.003 ± 0.007	[1.1, 2.5]	0.227 ± 0.023
[2.5, 4]	-0.12 ± 0.13	[2.5, 4]	-0.005 ± 0.017	[2.5, 4]	0.370 ± 0.021
[4, 6]	-0.50 ± 0.11	[4, 6]	-0.002 ± 0.017	[4, 6]	0.314 ± 0.046
[6, 8]	-0.73 ± 0.12	[6, 8]	$+0.002 \pm 0.010$	[6, 8]	0.216 ± 0.061

BUT, it only matters when discussing the **type** of NP we can see.