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# Real Scalar Dark Matter

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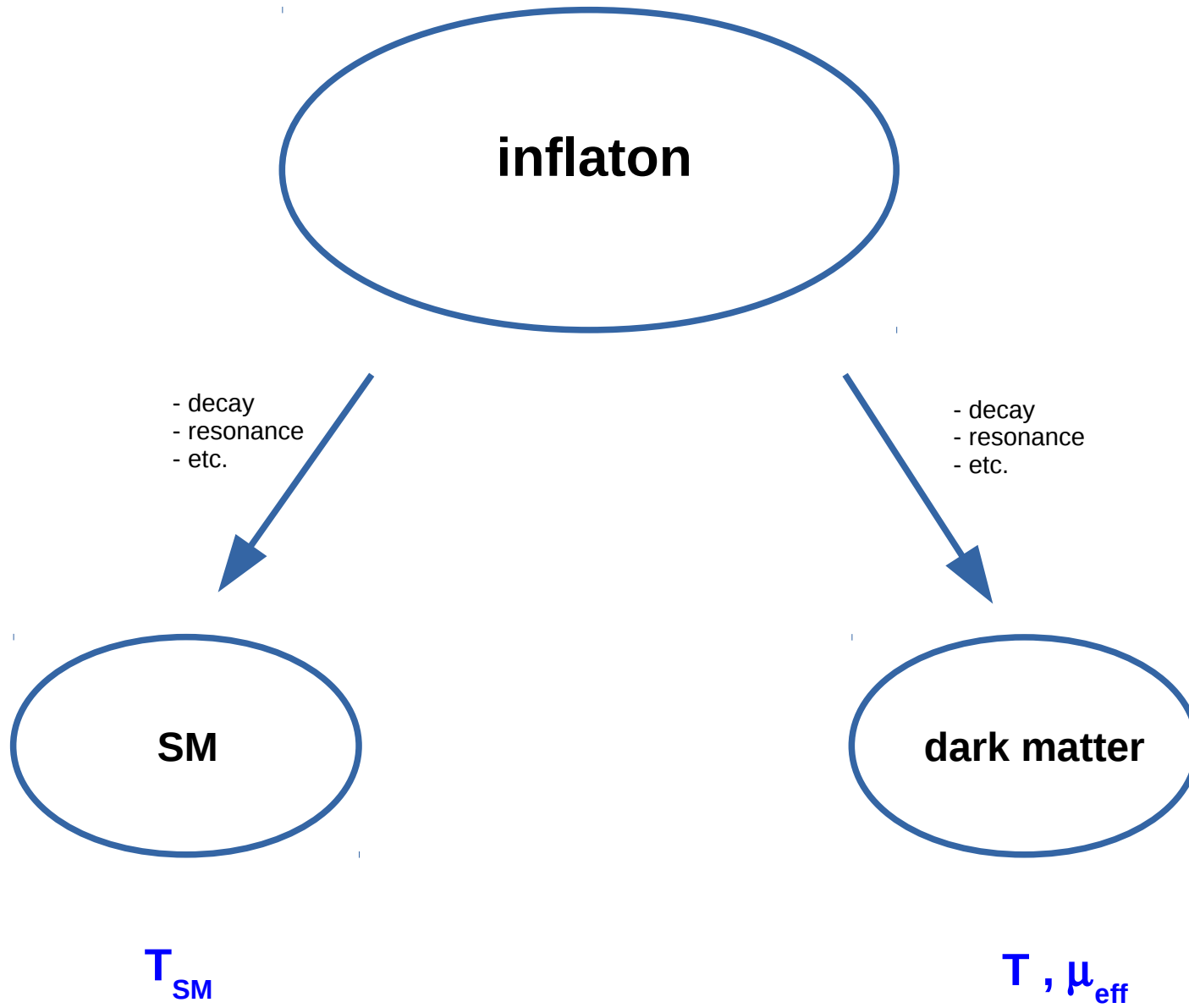
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## Based on

- 1906.07659 Arcadi, Lebedev, Pokorski, Toma  
“Real Scalar Dark Matter: Relativistic Treatment”
- 1908.05491 Lebedev, Toma  
“Relativistic Freeze-in”



# Simplest dark matter model

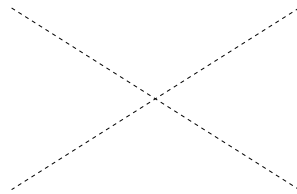
$$V = \frac{m^2}{2} S^2 + \frac{\lambda}{4!} S^4$$

Thermodynamics:

*Bose-Einstein statistics*

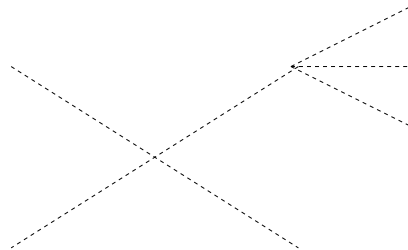
$$f(p) = \frac{1}{e^{\frac{E-\mu}{T}} - 1}$$

2 ↔ 2



$\Gamma_{22} > H \quad \rightarrow \quad T, \mu_{\text{eff}}$

2 ↔ 4



$\Gamma_{24} > H \quad \rightarrow \quad T$

## Dynamics:

initial state  $\rightarrow$  thermalization  $\rightarrow$  freeze-out

Common approach:

$$\left\{ \begin{array}{l} f(p) = \frac{1}{e^{\frac{E-\mu}{T}} - 1} \\ \mu \rightarrow 0 \end{array} \right. \rightarrow f(p) = e^{-(E-\mu)/T}$$

**Does not work in the relativistic regime  $T \gg m$  !**

## Relativistic reaction rates

$$\Gamma_{a \rightarrow b} = \int \left( \prod_{i \in a} \frac{d^3 \mathbf{p}_i}{(2\pi)^3 2E_i} f(p_i) \right) \left( \prod_{j \in b} \frac{d^3 \mathbf{p}_j}{(2\pi)^3 2E_j} (1 + f(p_j)) \right) |\mathcal{M}_{a \rightarrow b}|^2 (2\pi)^4 \delta^4(p_a - p_b)$$

2 → 4 :

$$\sigma(p_1, p_2) = \frac{1}{4F(p_1, p_2)} \int |\mathcal{M}_{2 \rightarrow 4}|^2 (2\pi)^4 \delta^4(p_1 + p_2 - \sum_i k_i) \prod_i \frac{d^3 \mathbf{k}_i}{(2\pi)^3 2E_{k_i}} (1 + f(k_i))$$

↑  
can't neglect!

$$\Gamma_{2 \rightarrow 4} = (2\pi)^{-6} \int d^3 \mathbf{p}_1 d^3 \mathbf{p}_2 f(p_1) f(p_2) \sigma(p_1, p_2) v_{M\emptyset}$$

**Issue:** can only compute  $\sigma(2 \rightarrow 4)$  in the center-of-mass frame (CalcHEP, etc.)

## Conversion to the CM frame:

$$p_1, p_2 \quad \rightarrow \quad p = \frac{p_1 + p_2}{2}, \quad k = \frac{p_1 - p_2}{2}$$

$$p = \Lambda(p) \begin{pmatrix} E \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

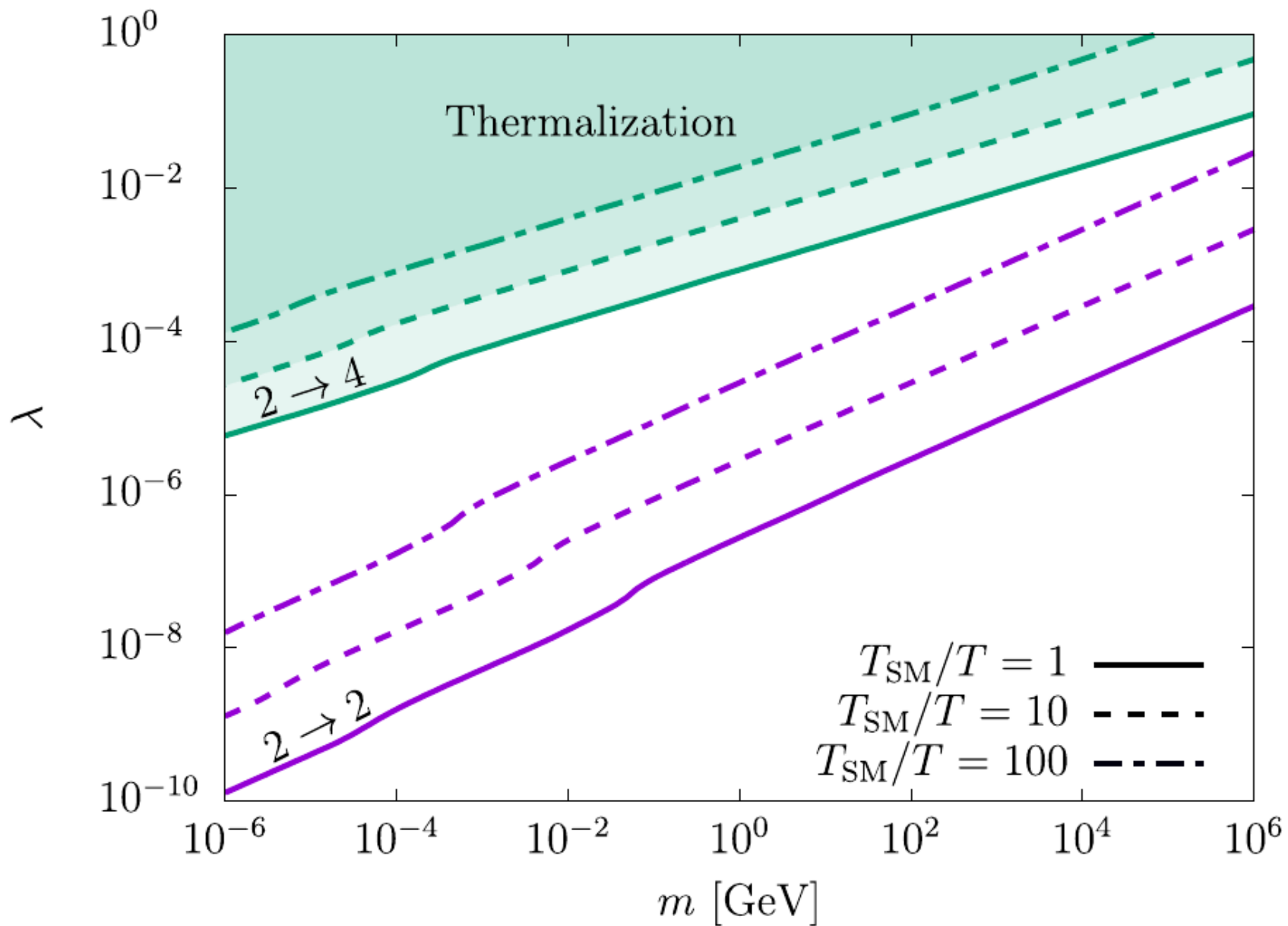
$$\rightarrow \int \frac{d^3 \mathbf{p}_1}{2E_1} \frac{d^3 \mathbf{p}_2}{2E_2} \dots = 2 \int_m^\infty dE \sqrt{E^2 - m^2} E^2 \int_0^\infty d\eta \sinh^2 \eta \int d\Omega_p d\Omega_k \dots$$

## Relativistic analog of the Gelmini-Gondolo formula:

$$\Gamma_{2 \rightarrow 4} = \frac{4T}{\pi^4} \int_m^\infty dE E^3 \sqrt{E^2 - m^2} \int_0^\infty d\eta \frac{\sinh \eta}{e^{2(E \cosh \eta - \mu)/T} - 1} \ln \frac{\sinh \frac{E \cosh \eta + \sqrt{E^2 - m^2} \sinh \eta - \mu}{2T}}{\sinh \frac{E \cosh \eta - \sqrt{E^2 - m^2} \sinh \eta - \mu}{2T}} \sigma_{\text{CM}}(E, \eta)$$

includes BE final state factors

Thermal or kinetic equilibrium:  $\Gamma_{24}$  or  $\Gamma_{22} > n H$





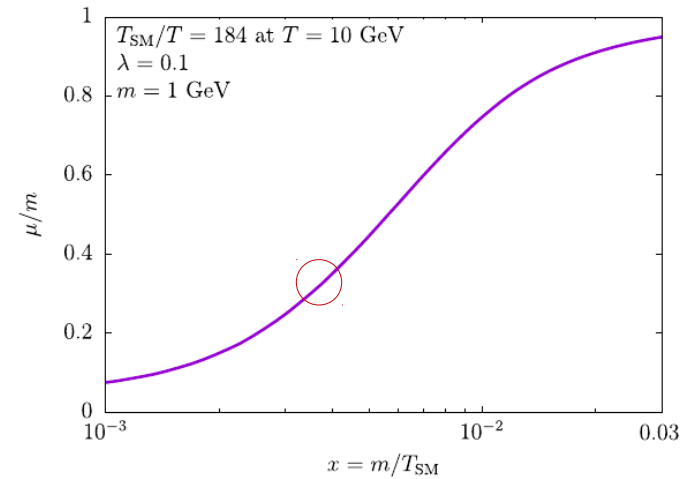
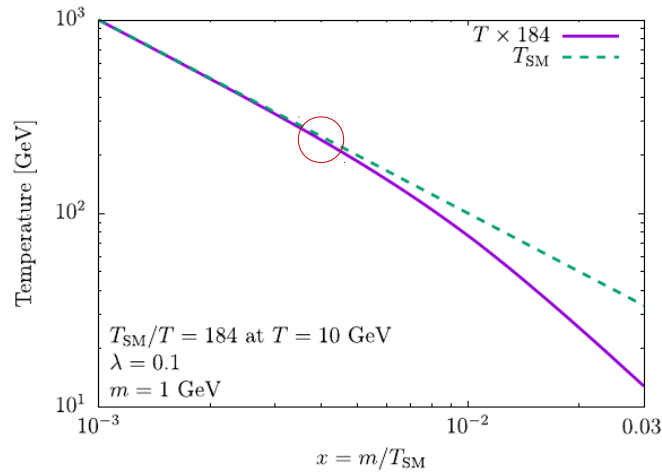
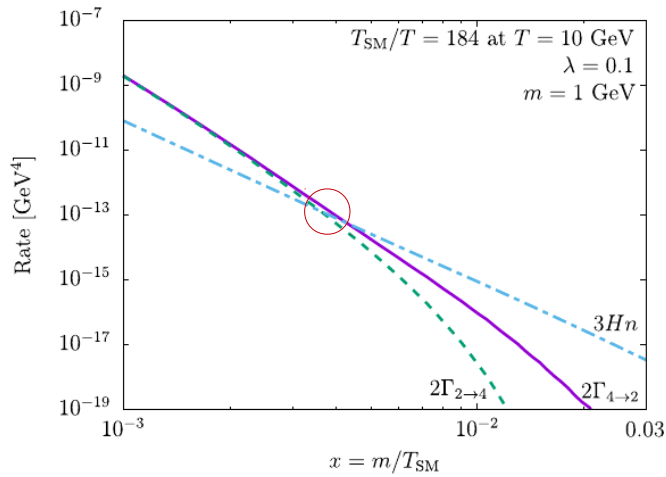
## Boltzmann equation + entropy conservation:

$$\left\{ \begin{array}{l} \frac{dn}{dt} + 3Hn = 2 (\Gamma_{2 \rightarrow 4} - \Gamma_{4 \rightarrow 2}) \\ \frac{s}{s_{\text{SM}}} = \text{const} \end{array} \right.$$



$T(t), \mu(t)$

## Freeze-out $2 \Gamma_{24} < 3 n H$ :



## Thermal mass effect at high T:

$$m^2 \rightarrow m^2 + \frac{\lambda}{24} T^2$$

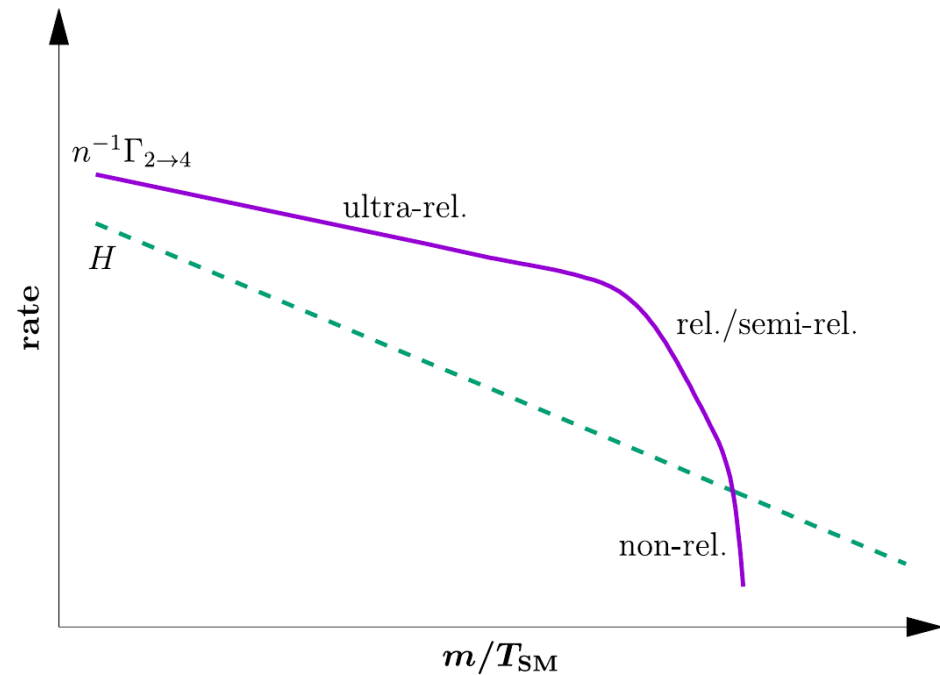
$$\Gamma_{2 \rightarrow 2} \propto T^4 \ln \frac{T}{m_{\text{eff}}} \rightarrow c T^4$$

IR divergent without  
thermal mass

## Freeze-out can't be ultra-relativistic:

$$\Gamma_{24} \sim T^4$$

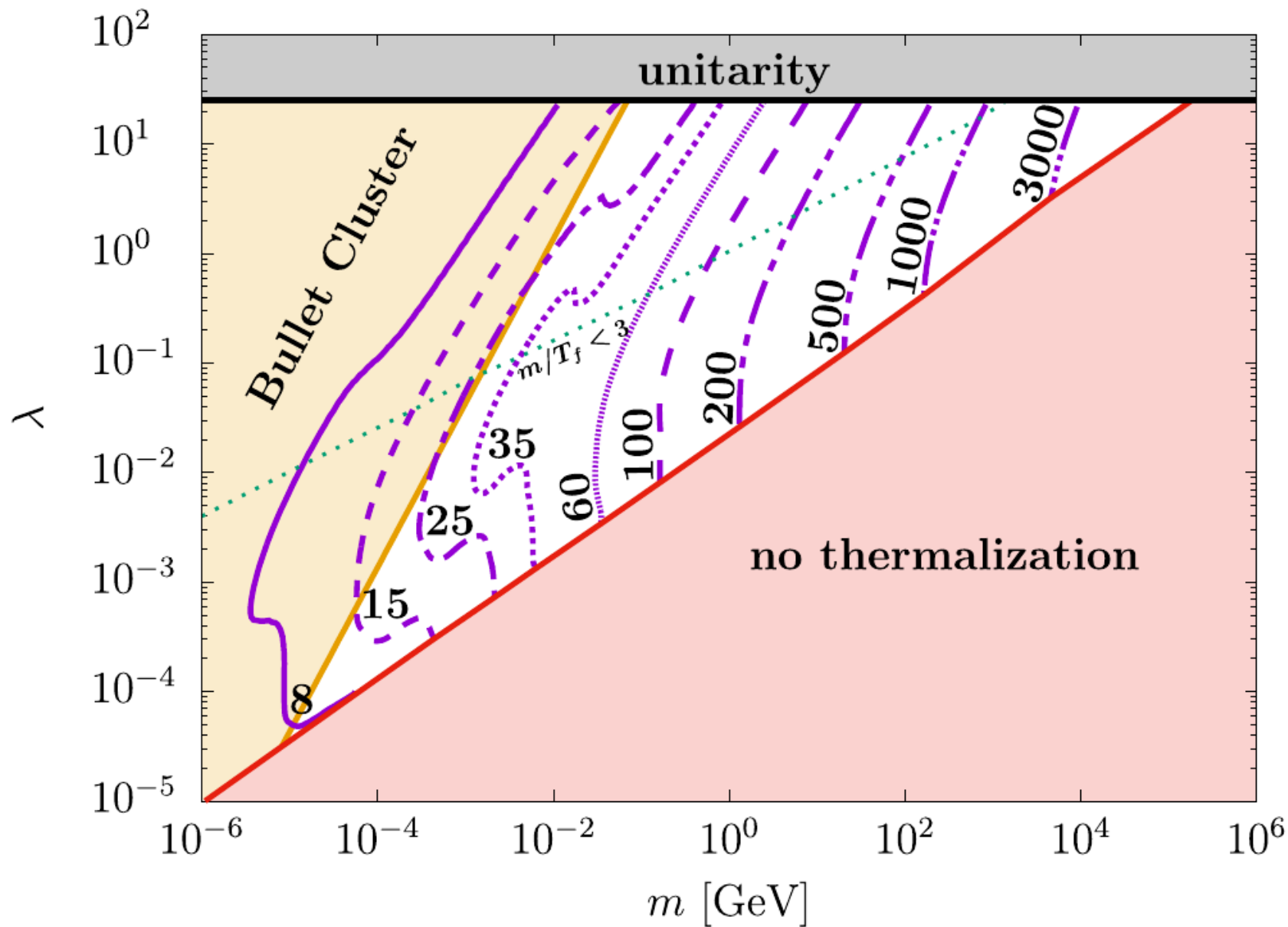
$$n H \sim T^5$$



# Correct DM relic abundance:

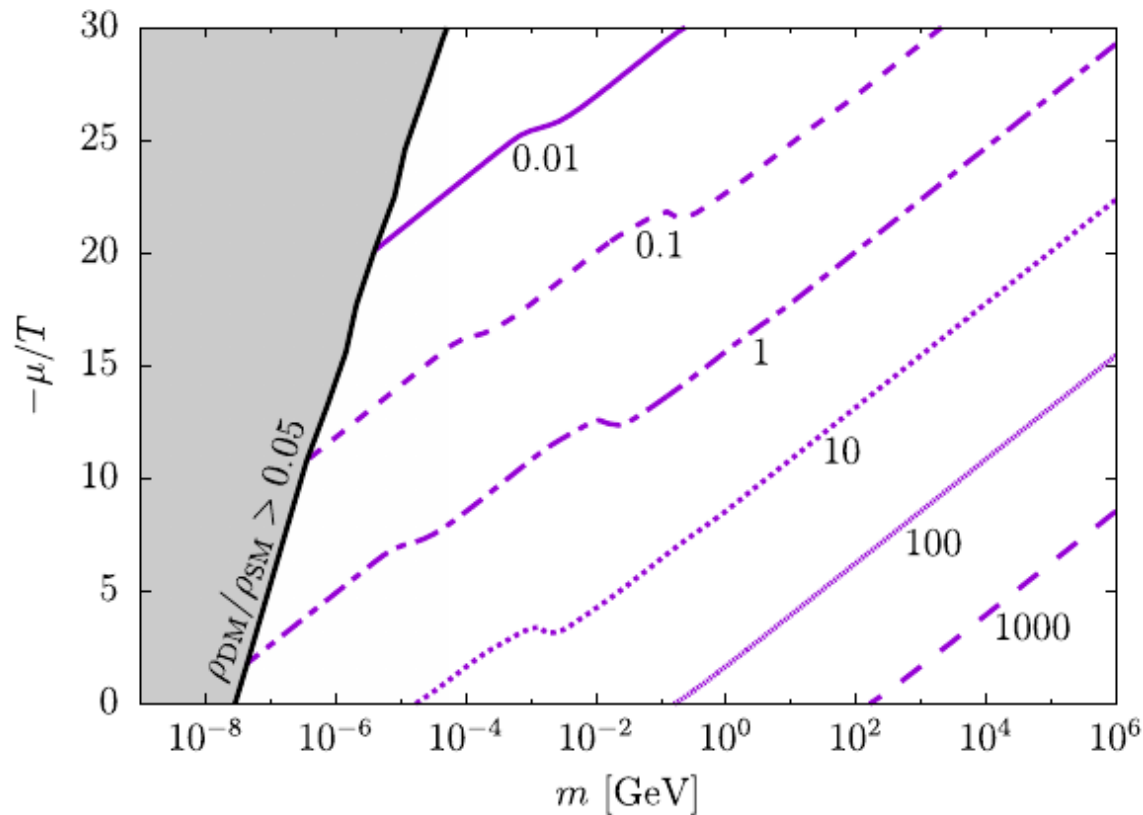
----- = boundary of relativistic freeze-out

8, 15, 25, ... =  $T_{SM} / T$



Only kinetic equilibrium:

$$n = \frac{T^3}{\pi^3} \text{Li}_3(e^{\mu/T}) \quad , \quad \mu \propto T \propto T_{\text{SM}}$$



white = correct relic density  
 0.01, 0.1, 1, ... =  $T_{\text{SM}}/T$

**NB:**  $|\mu| < m$  if there are “non-trivial” anti-particles (Haber, Weldon ‘81)

# Relativistic freeze-in

$$V_{hs} = \frac{1}{2} \lambda_{hs} H^\dagger H s^2$$

$T > T_{EW}$  regime:

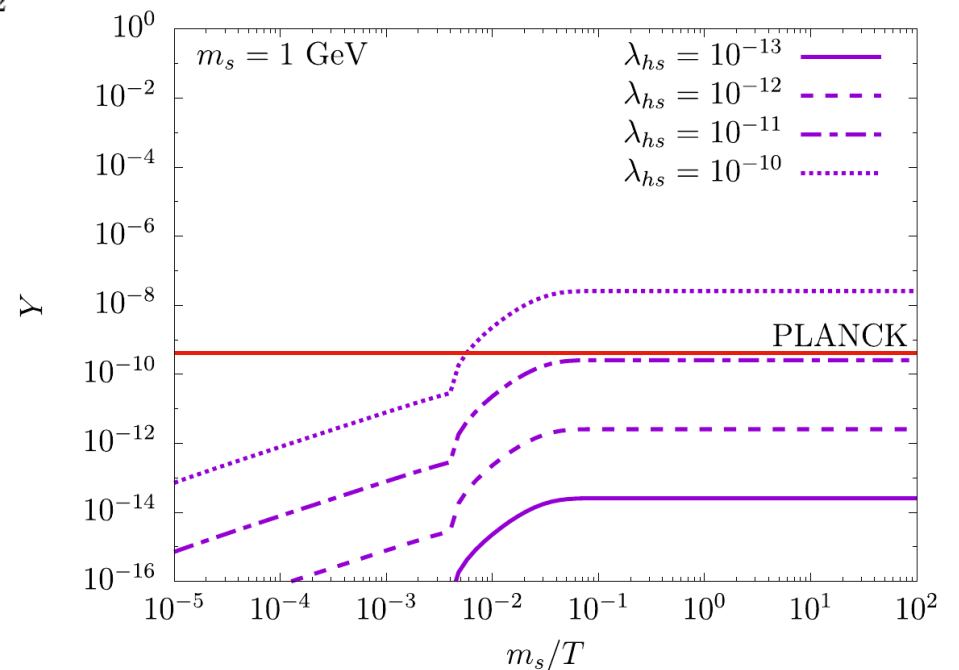
$h_i h_i \rightarrow ss$

$$\Gamma_{2 \rightarrow 2} = \frac{1}{2!2!} \frac{\lambda_{hs}^2 T}{16\pi^5} \int_{m_h}^{\infty} dE E \sqrt{E^2 - m_s^2} \int_0^{\infty} d\eta \frac{\sinh \eta}{e^{\frac{2E}{T} \cosh \eta} - 1} \ln \frac{\sinh \frac{E \cosh \eta + \sqrt{E^2 - m_h^2} \sinh \eta}{2T}}{\sinh \frac{E \cosh \eta - \sqrt{E^2 - m_h^2} \sinh \eta}{2T}}$$

$$m_h^2 \simeq m_{h0}^2 + \left( \frac{3}{16} g_2^2 + \frac{1}{16} g_1^2 + \frac{1}{4} y_t^2 + \frac{1}{2} \lambda_h \right) T^2$$

$$m_s > m_h \quad \rightarrow \quad \lambda_{hs} \simeq 2.2 \times 10^{-11}$$

$$m_s \ll m_h \quad \rightarrow \quad \lambda_{hs} \simeq 1.2 \times 10^{-11} \sqrt{\frac{\text{GeV}}{m_s}}$$



## Conclusion:

- scalar DM evolution in the relativistic regime
- Bose-Einstein vs Maxwell-Boltzmann
- effective chemical potential
- relativistic effects in freeze-in