

On asymptotic symmetries in spatially flat FLRW

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Early stage:

- Asymptotic structure of asymptotically flat spacetimes at \mathcal{I}^+ [Bondi, van der Burg & Metzner (1962) and Sachs (1962)].
- AdS₃/CFT₂ correspondence [Brown & Henneaux (1986)] → AdS/CFT correspondence [Maldacena (1999)].

Modern era:

- Inclusion of superrotations [de Boer & Solodukhin (2003) and Barnich & Tossaert (2010)].
- Infrared triangle: asymptotic symmetries ↔ memory effects ↔ soft theorems [Strominger & Zhiboedov (2014) and Strominger (2017)].

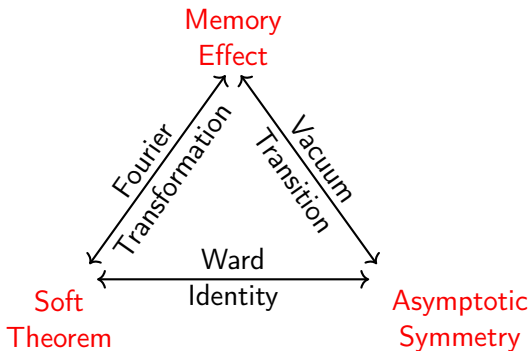


Figure: The infrared triangle according to [Strominger (2017)]

Introduction - Objectives and applications

Diversification into a variety of topics, including flat holography [Strominger (2014), Stieberger & Taylor (2019), Donnay et al. (2020)], black hole entropy [Hawking, Perry & Strominger (2016), Averin, Gómez, Dvali & Lüst (2016), Donnay et al. (2016), Grumiller et al. (2020)], extension to spatial and timelike flat infinity [Prabhu et al. (2019), Khairnar et al. (2000)], extension to (A)dS₄ [Compère et al. (2000)], string theory [Avery et al. (2016), Afshar et al. (2019)], swampland [Lüst et al. (2020)], membrane paradigm/fluid-gravity duality [Penna (2016), Penna (2017)] ...

... **but what about cosmological spacetimes?** Very little has been done.

- Cosmological gravitational memory effect [Kehagias et al. (2016), Wald et al. (2016), Shiu et al. (2017)].
- Asymptotic diffeomorphisms in cosmological spacetimes [Shiromizu et al. (1999), Ferreira et al. (2017), Bonga & Prabhu (2020)].

Spatially flat FRW spacetimes are conformal to flat spacetimes. In Bondi coordinates

$$u = \eta - \sqrt{x^i x_i}, \quad r = \sqrt{x^i x_i}, \quad z = \frac{x^1 + ix^2}{x^3 + \sqrt{x^i x_i}}, \quad \bar{z} = \frac{x^1 - ix^2}{x^3 + \sqrt{x^i x_i}},$$

their metrics can be written as

$$ds^2 = \left(\frac{r+u}{L} \right)^{2k} \left(-du^2 - 2dudr + \frac{4r^2}{(1+z\bar{z})^2} dzd\bar{z} \right),$$

where η is the conformal time and $k = 2/(3\omega + 1)$, $\omega = p/\rho$.

These spacetimes can be divided into accelerated ($-1 < \omega < -1/3$, $k < 0$) and decelerated ($-1/3 < \omega < 1$, $k > 0$) expansion.

It turns out that only the decelerated have future null infinity \mathcal{I}^+ → analysis restricts to them.

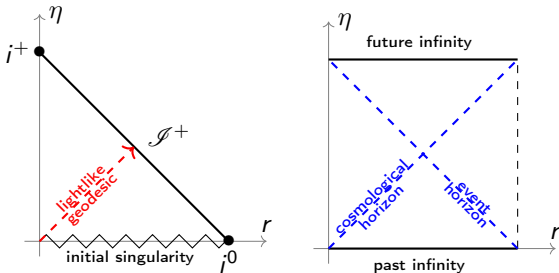


Figure: Conformal diagram of spatially flat decelerating FLRW (left) and accelerating FLRW (right)

Our project: analysis of asymptotically decelerating spatially flat FLRW metrics at \mathcal{I}^+ . Concretely, we delve into the *asymptotic symmetry corner of the cosmological infrared triangle*, aiming to extend the flat space analysis to a more phenomenologically relevant framework.

This is an unavoidable step towards applicability and testability of the theory.

Requirements:

- Vacuum space is unperturbed FLRW.
- Compatible with Bondi gauge and frame.
- Allow for general cosmological perturbations but preserve (to leading order) homogeneous, isotropic and spatially flat profile, as well as k -fluid background.
- Closed in the r -expansion upon application of asymptotic diffeomorphism.
- Trace and components of the energy momentum tensor remains at most finite when appropriately scaled.

Ansatz:

$$ds^2 = \left(\frac{r+u}{L}\right)^{2k} \left\{ - \left(1 - \Phi - \frac{2m}{r}\right) du^2 - 2 \left(1 - \Psi - \frac{K}{r}\right) dudr - 2(r\Theta_A + U_A + \frac{1}{r}N_A)dudx^A + ((1 + \Omega)r^2\gamma_{AB} + rC_{AB} + h_{AB}) dx^A dx^B \right\} .$$

(Some) Results

Asymptotic diffeomorphisms:

$$\begin{aligned}\xi = & \xi^u(u, z, \bar{z})\partial_u + \left[r\xi^{r(V)}(z, \bar{z}) + \xi^{r(0)} + \frac{1}{r}\xi^{r(1)} \right] \partial_r \\ & + \left[V^B(z, \bar{z}) + \frac{1}{r}\xi^{B(1)} + \frac{1}{r^2}\xi^{B(2)} \right] \partial_B ,\end{aligned}$$

with

$$\xi^{B(1)} = -\frac{(1-\Psi)}{1+\Omega} D^B \xi^u ,$$

$$\xi^{B(2)} = \frac{1}{2} \left(\frac{1-\Psi}{(1+\Omega)^2} C^{AB} D_A \xi^u + \frac{K}{1+\Omega} D^B \xi^u \right)$$

$$\xi^{r(0)} = \frac{1}{1+k} \left[-\frac{1}{2(1+\Omega)} D_A \left((1+\Omega) \xi^{A(1)} \right) - \frac{1}{2(1+\Omega)} \Theta^A D_A \xi^u + k u \xi^{r(V)} - k \xi^u \right]$$

$$\begin{aligned}\xi^{r(1)} = & \frac{1}{2(1+k)(1+\Omega)} \left[\frac{C_B^A \Theta_A D^B \xi^u}{1+\Omega} - 2k(1+\Omega) \left(u^2 \xi^{r(V)} - u \xi^{r(0)} - u \xi^u \right) \right. \\ & \left. - D_A \left((1+\Omega) \xi^{A(2)} \right) - U^A D_A \xi^u \right] .\end{aligned}$$

Effect on asymptotic data:

$$\delta\Phi = V^A D_A \Phi + \xi^u \partial_u \Phi - 2(1 - \Psi) \partial_u \xi^{r(0)} - 2k(1 - \Phi) \xi^{r(V)} - 2(1 - \Phi) \partial_u \xi^u + 2\Theta_A \partial_u \xi^{A(1)}$$

$$\delta m = \xi^u \partial_u m - k(1 - \Phi) \xi^u - ((1 - 2k)m - ku(1 - \Phi)) \xi^{r(V)} - k(1 - \Phi) \xi^{r(0)} + V^A D_A m + \frac{1}{2} \xi^{A(1)} D_A \Phi + K \partial_u \xi^{r(0)} - (1 - \Psi) \partial_u \xi^{r(1)} + m \partial_u \xi^u + U_A \partial_u \xi^{A(1)} + \Theta_A \partial_u \xi^{A(2)}$$

$$\delta\Psi = V^A \partial_A \Psi + \xi^u \partial_u \Psi - (1 + 2k)(1 - \Psi) \xi^{r(V)} - (1 - \Psi) \partial_u \xi^u$$

$$\delta K = \xi^u \partial_u K + V^A D_A K + K \partial_u \xi^u + \xi^{A(1)} D_A \Psi - \Theta_A \xi^{A(1)} + 2k(1 - \Psi) \left(u \xi^{r(V)} - \xi^u - \xi^{r(0)} \right) + 2kK \xi^{r(V)}$$

$$\delta\Omega = V^C D_C \Omega + \xi^u \partial_u \Omega + 2(1 + k) \xi^{r(V)} + (1 + \Omega) D_A V^A$$

$$\begin{aligned}
\delta C_{AB} = & \xi^u \partial_u C_{AB} + V^C D_C C_{AB} + C_{AC} D_B V^C + C_{BC} D_A V^C \\
& + 2(1 + \Omega) \gamma_{AB} ((1 + k) \xi^{r(0)} - k u \xi^{r(V)} + k \xi^u) + \gamma_{AB} \xi^{C(1)} D_C \Omega \\
& + (1 + \Omega) (D_A \xi_B^{(1)} + D_B \xi_A^{(1)}) + \Theta_A D_B \xi^u + \Theta_B D_A \xi^u \\
& + (1 + 2k) C_{AB} \xi^{r(V)}
\end{aligned}$$

$$\begin{aligned}
\delta \Theta_A = & V^B D_B \Theta_A + \xi^u \partial_u \Theta_A + (1 + 2k) \Theta_A \xi^{r(V)} + \Theta_B D_A V^B \\
& - (1 - \Psi) \partial_A \xi^{r(V)} + \Theta_A \partial_u \xi^u + (1 + \Omega) \partial_u \xi_A^{(1)}
\end{aligned}$$

$$\begin{aligned}
\delta U_A = & (2k \Theta_A + \partial_u U_A) \xi^u + (1 + 2k) \Theta_A \xi^{r(0)} + 2k \xi^{r(V)} (U_A - u \Theta_A) \\
& + V^B D_B U_A + \xi^{B(1)} D_B \Theta_A + \Theta_B D_A \xi^{B(1)} + U_B D_A V^B \\
& - (1 - \Psi) D_A \xi^{r(0)} + K D_A \xi^{r(V)} - (1 - \Phi) D_A \xi^u + U_A \partial_u \xi^u \\
& + C_{AB} \partial_u \xi^{B(1)} + (1 + \Omega) \partial_u \xi_A^{(2)}
\end{aligned}$$

- Introduced the class of metrics to be considered asymptotically decelerating spatially flat FLRW at future null infinity \mathcal{I}^+ .
- Obtained supertranslation and superrotation-like diffeomorphisms together with their effect on the asymptotic data.
- Investigated how these transformations act on a pure FRW background and noticed that, not only C_{zz} and $C_{\bar{z}\bar{z}}$, but all the asymptotic data can be generated out of the vacuum.
- Furthermore, we have realized that central inhomogeneities are more involved to describe than Schwarzschild in the flat case. The Sultana-Dyer black hole background, has been explored and we have realized that it is not covered by our ansatz.
- Considered simple cosmologically perturbed backgrounds and analyzed how the perturbations affect the remaining asymptotic data under the action of the asymptotic diffeomorphisms.

Current and future research

Working at present:

- Extend our ansatz to include anisotropic spacetimes (e.g. Bianchi Type I).
- Extend our ansatz to include cosmological black holes (e.g. Sultana-Dyer and primordial black holes).
- Understanding the asymptotic symmetry algebra.
- Cosmological holographic flow.

Long-term goals:

- On-shell implementation for GR and alternative gravity theories.
- Extension to accelerating spatially flat FLRW spacetimes.*
- Entropy of cosmological black holes (and horizons).
- Cosmological (non-linear) memory effects (scalar, vector and tensor).
- Scattering problem in cosmology.

Thank you for your attention!