



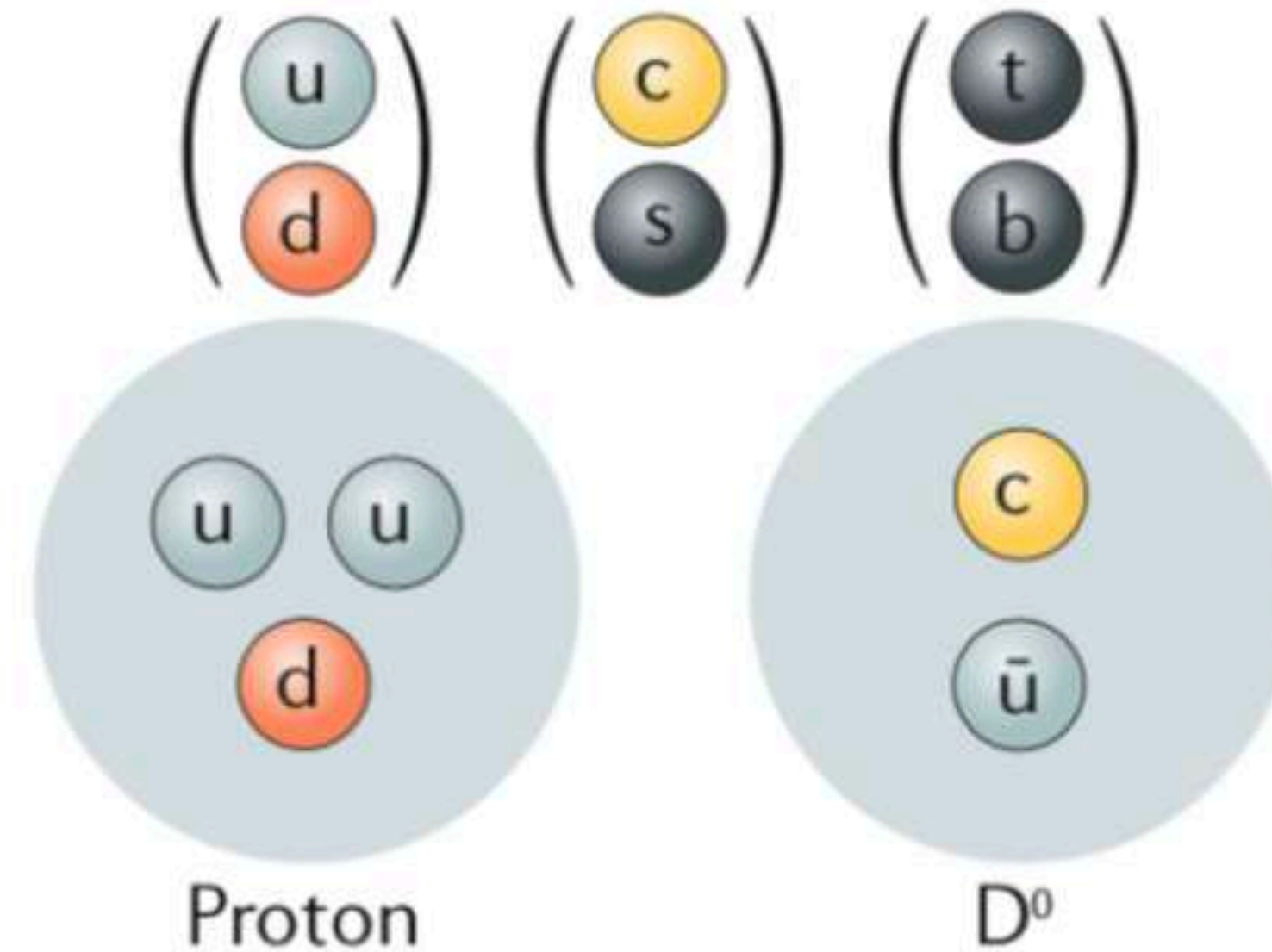
SM predictions for D **Mixing**

MIAPP

22.3.2022

Alexander Lenz

Charm Physics



	$D^0 = (\bar{u}c)$	$D^+ = (\bar{d}c)$	$D_s^+ = (\bar{s}c)$	$\Lambda_c = (udc)$
Mass (GeV)	1.86486	1.86962	1.96850	2.28646
Lifetime (ps)	0.4101	1.040	0.500	0.200

Theoretical Peculiarities of Charm:

1. The strong coupling is strong

$$\alpha_s(m_c) = 0.33 \pm 0.01$$

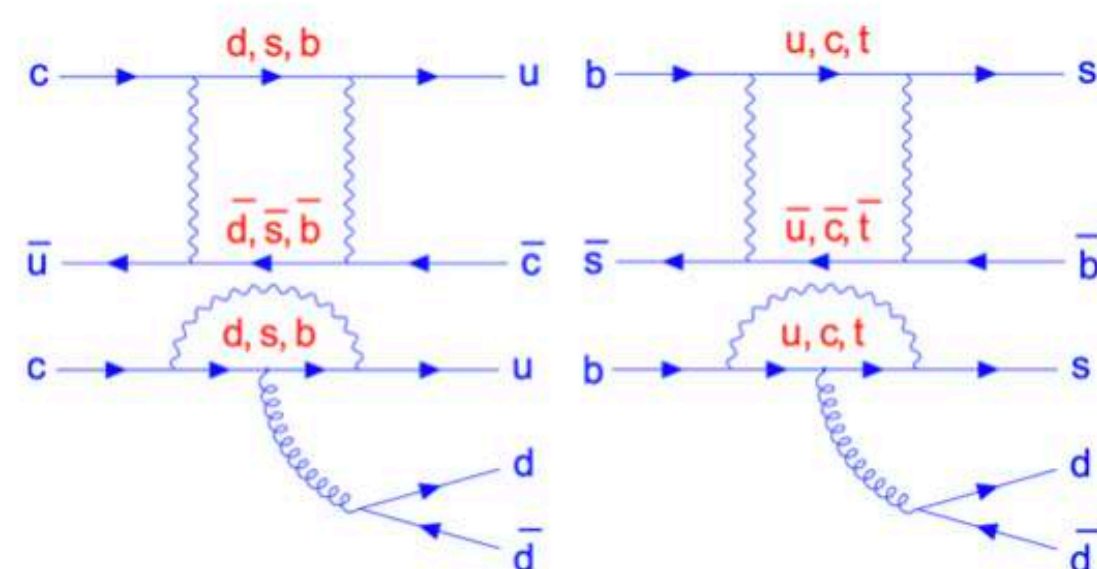
2. The charm quark is not really heavy

$$m_c^{\text{Pole}} = (1.67 \pm 0.07) \text{ GeV}, \quad \bar{m}_c(\bar{m}_c) = (1.27 \pm 0.02) \text{ GeV},$$

3. There is almost no CPV in charm

$$V_{cd} = -0.2247 - 1.4 \cdot 10^{-4}I, \quad V_{cs} = 0.97354 - 3.1 \cdot 10^{-5}I, \quad V_{cb} = 0.0416$$

4. There are extremely pronounced GIM cancellations in the charm sector



$$\begin{aligned} \left(\frac{m_d}{M_W}\right)^2 &\approx 0, & \left(\frac{m_u}{M_W}\right)^2 &\approx 0, \\ \left(\frac{m_s}{M_W}\right)^2 &\approx 1.3 \cdot 10^{-6}, & \left(\frac{m_c}{M_W}\right)^2 &\approx 2.5 \cdot 10^{-4}, \\ \left(\frac{m_b}{M_W}\right)^2 &\approx 2.8 \cdot 10^{-3}, & \left(\frac{m_t}{M_W}\right)^2 &\approx 4.5. \end{aligned}$$

See e.g.
AL, G. Wilkinson
2011.04443



Thanks Tommaso and Patricia!

The charm system is theoretically more difficult than the b system since

$$\alpha_s(m_c) \approx 0.33 \quad \text{and} \quad \frac{\Lambda_{QCD}}{m_c} \approx 3 \frac{\Lambda_{QCD}}{m_b}$$

Nevertheless the **Heavy Quark Expansion** might still converge

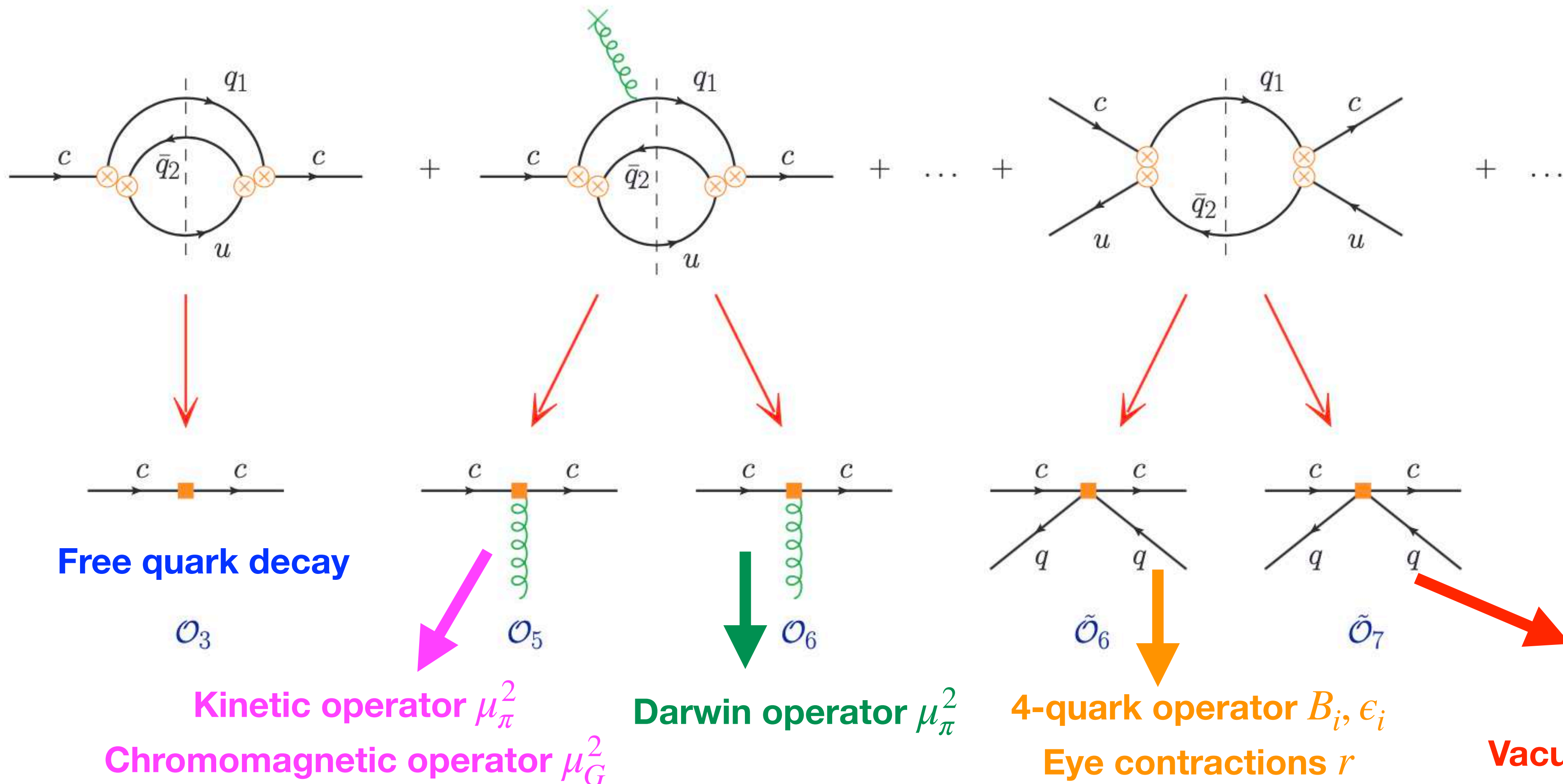
But things will become very ugly, if in addition cancellations arising

- A. No cancellations, e.g. $\Gamma(D^0)$
- B. Strong cancellations, e.g. $\Gamma(D^+)$
- C. Crazy cancellations, e.g. D -mixing

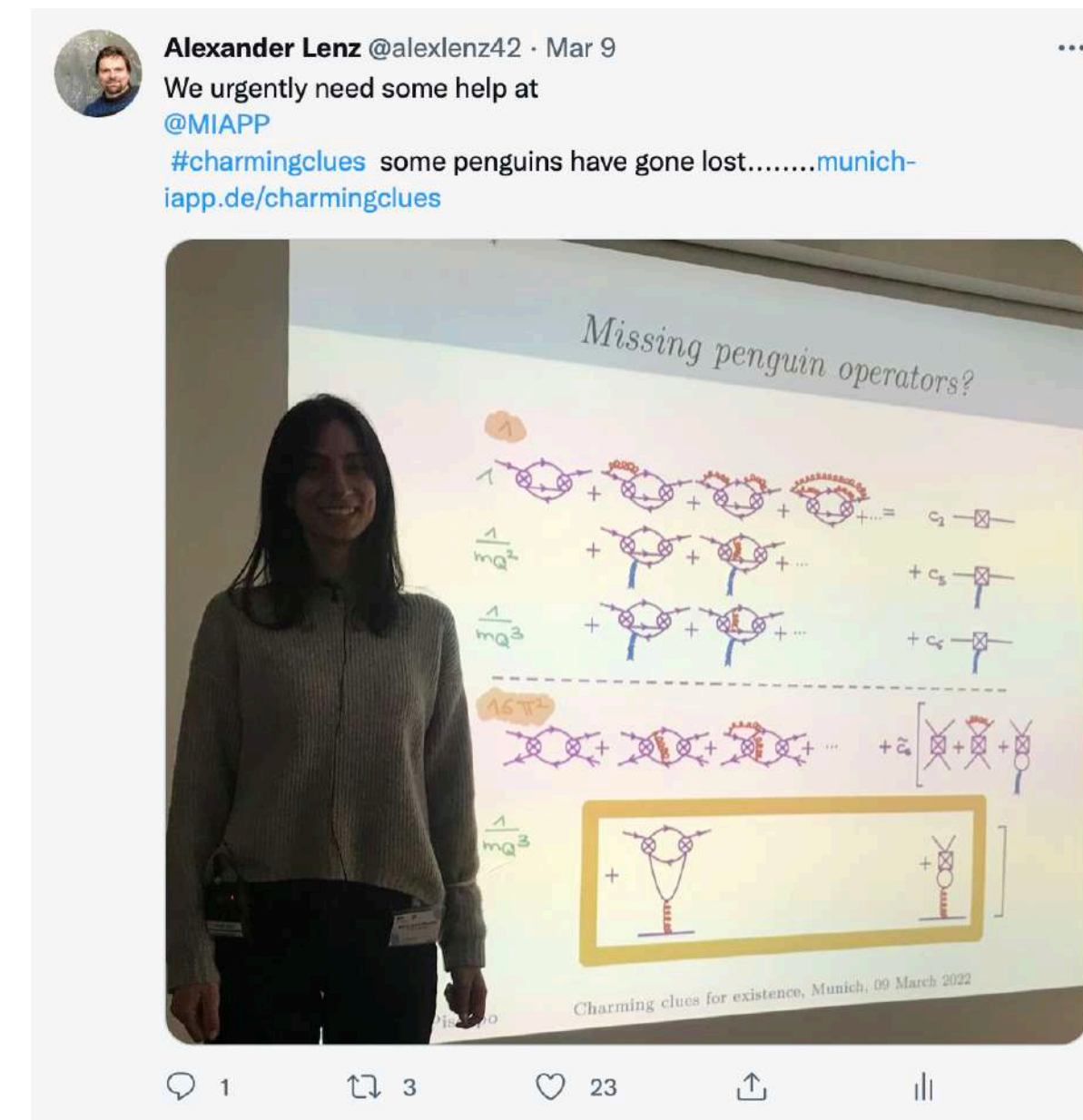


A. No Cancellations

$$\Gamma(D) = \Gamma_3 + \Gamma_5 \frac{\langle \mathcal{O}_5 \rangle}{m_c^2} + \Gamma_6 \frac{\langle \mathcal{O}_6 \rangle}{m_c^3} + \dots + 16\pi^2 \left(\tilde{\Gamma}_6 \frac{\langle \tilde{\mathcal{O}}_6 \rangle}{m_c^3} + \tilde{\Gamma}_7 \frac{\langle \tilde{\mathcal{O}}_7 \rangle}{m_c^4} + \dots \right), \quad \Gamma_i = \Gamma_i^{(0)} + \frac{\alpha_s(m_c)}{4\pi} \Gamma_i^{(1)} + \left[\frac{\alpha_s(m_c)}{4\pi} \right]^2 \Gamma_i^{(2)} + \dots$$



For more details, see the talk of Maria Laura

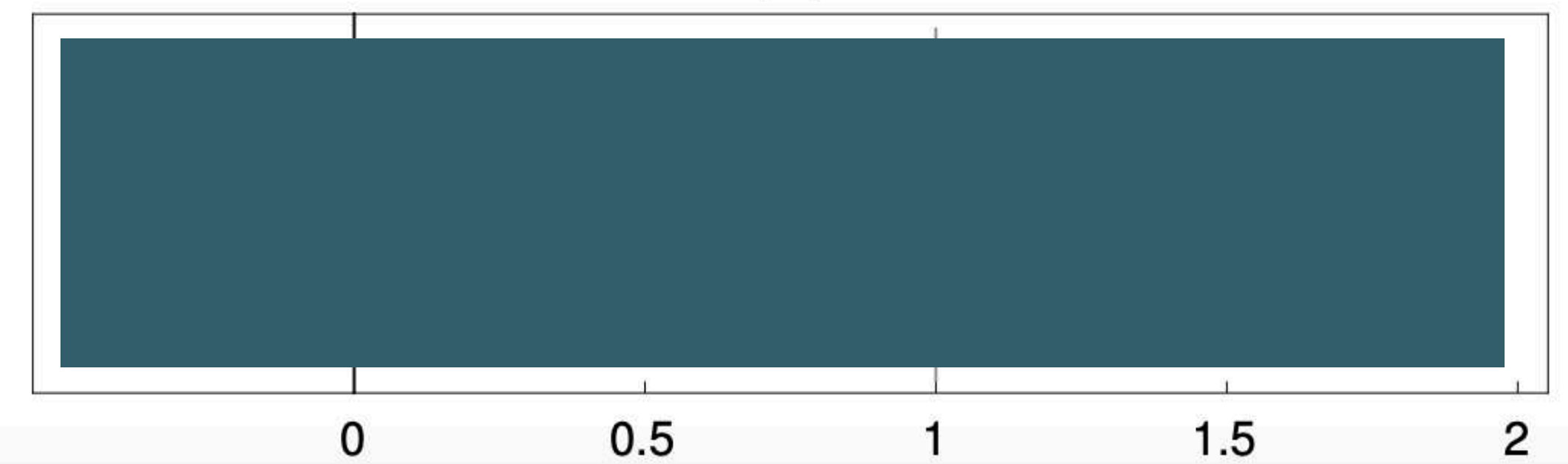
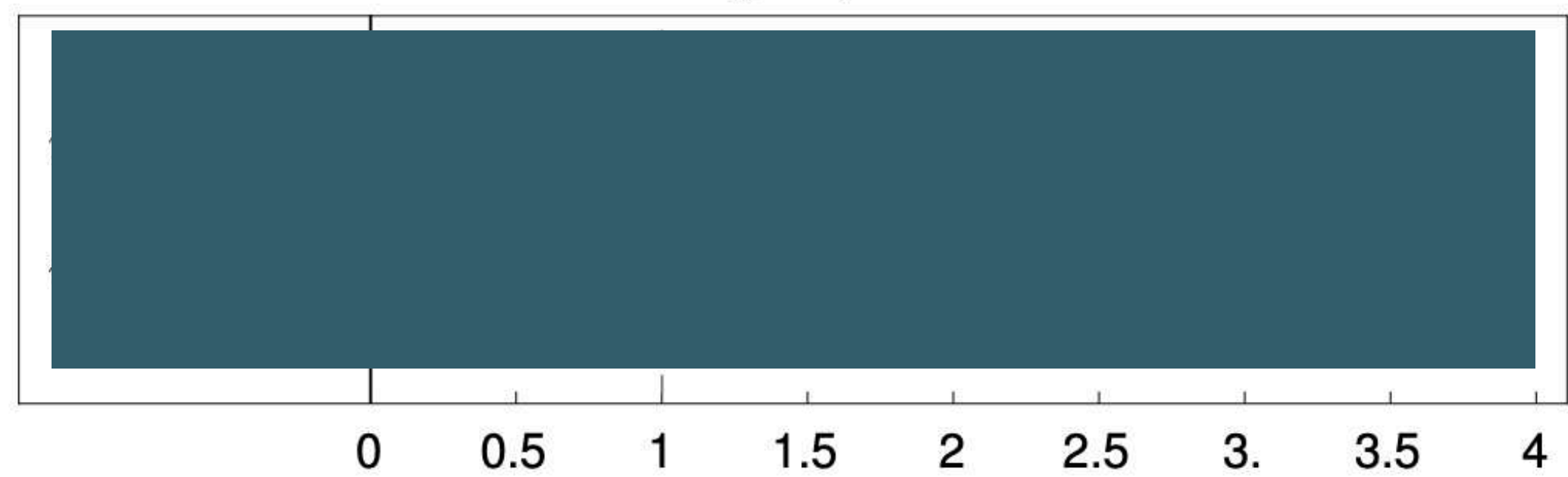
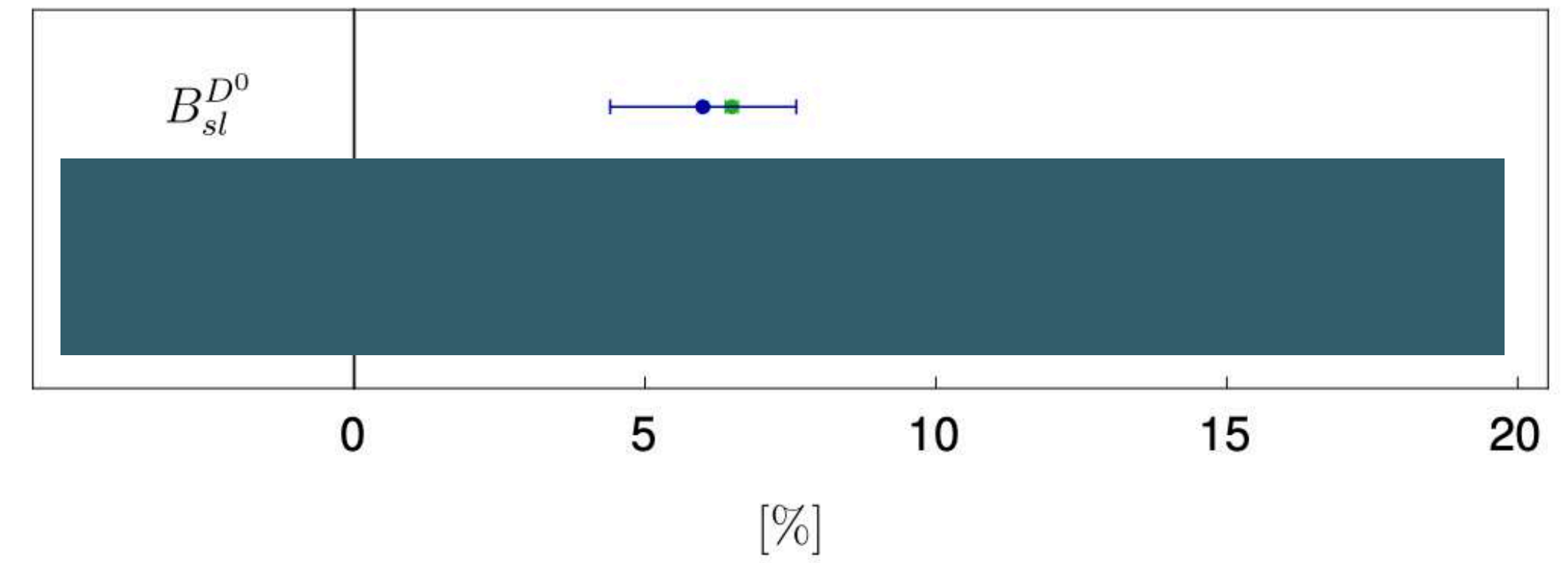
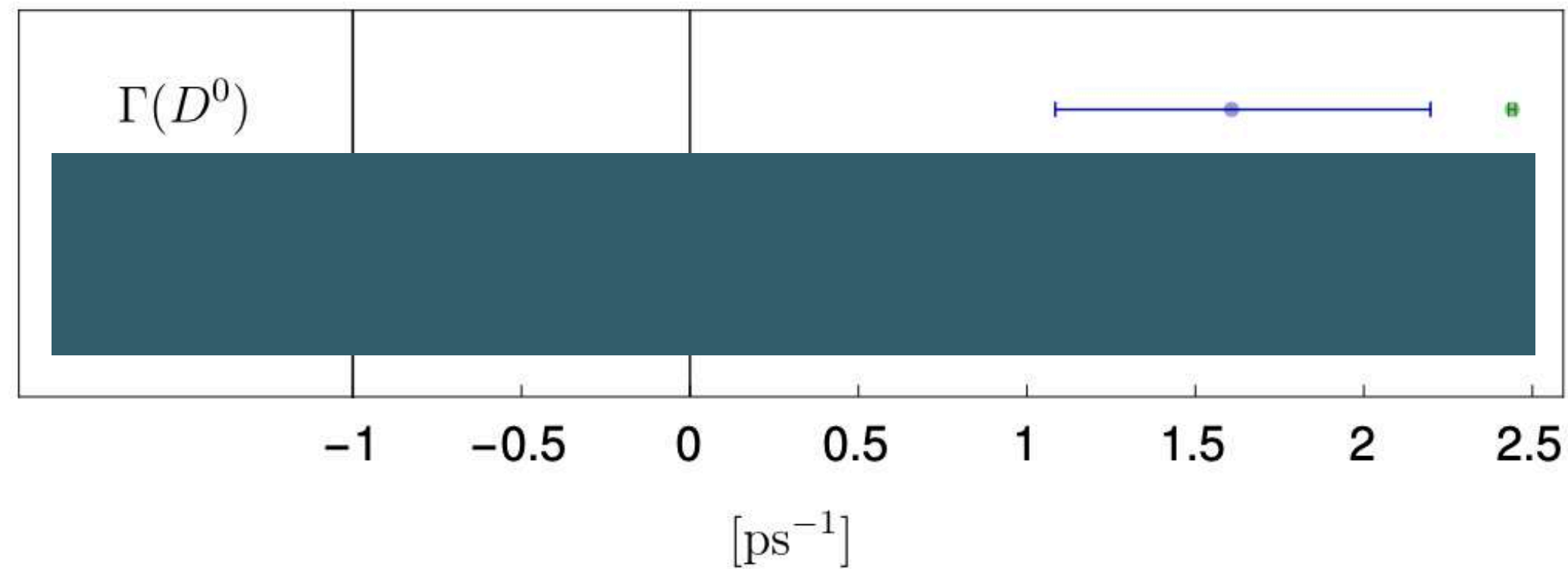


A. No Cancellations

$$\Gamma(D^0) = 6.15 \Gamma_0 \left[1 + 0.48 - 0.13 \frac{\mu_\pi^2(D)}{0.48 \text{ GeV}^2} + 0.01 \frac{\mu_G^2(D)}{0.34 \text{ GeV}^2} + 0.31 \frac{\rho_D^3(D)}{0.082 \text{ GeV}^3} \right.$$

$$\begin{aligned} & - \underbrace{0.01}_{\text{dim-6, VIA}} - 0.005 \frac{\delta \tilde{B}_1^q}{0.02} + 0.005 \frac{\delta \tilde{B}_2^q}{0.02} + 0.137 \frac{\tilde{\epsilon}_1^q}{-0.04} - 0.125 \frac{\tilde{\epsilon}_2^q}{-0.04} + \underbrace{0.00}_{\text{dim-7, VIA}} \\ & - 0.0045 r_1^{qq} - 0.0004 r_2^{qq} - 0.0035 r_3^{qq} + 0.0000 r_4^{qq} \\ & - 0.0109 r_1^{sq} - 0.0079 r_2^{sq} - 0.0000 r_3^{sq} + 0.0001 r_4^{sq} \end{aligned} \left. \right].$$

A. No Cancellations



- Values of $\mu_\pi^2, \mu_G^2, \rho_D^3$ almost unknown
- NNLO-QCD corrections to free quark decay in progress
Fael, Steinhauser,...
- NNLO-QCD corrections to spectator effects in progress
Nierste, Steinhauser,...

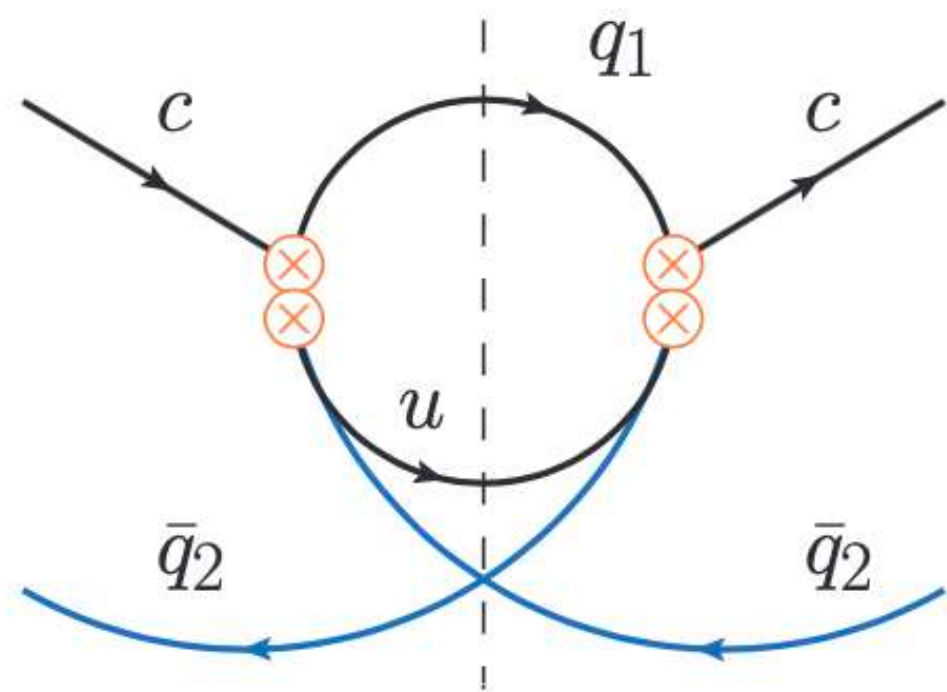


Revisiting Inclusive Decay Widths of Charmed Mesons #5

Daniel King (Durham U., IPPP and Durham U.), Alexander Lenz (Siegen U.), Maria Laura Piscopo (Siegen U.), Thomas Rauh (U. Bern, AEC), Aleksey V. Rusov (Siegen U.) et al. (Sep 27, 2021)

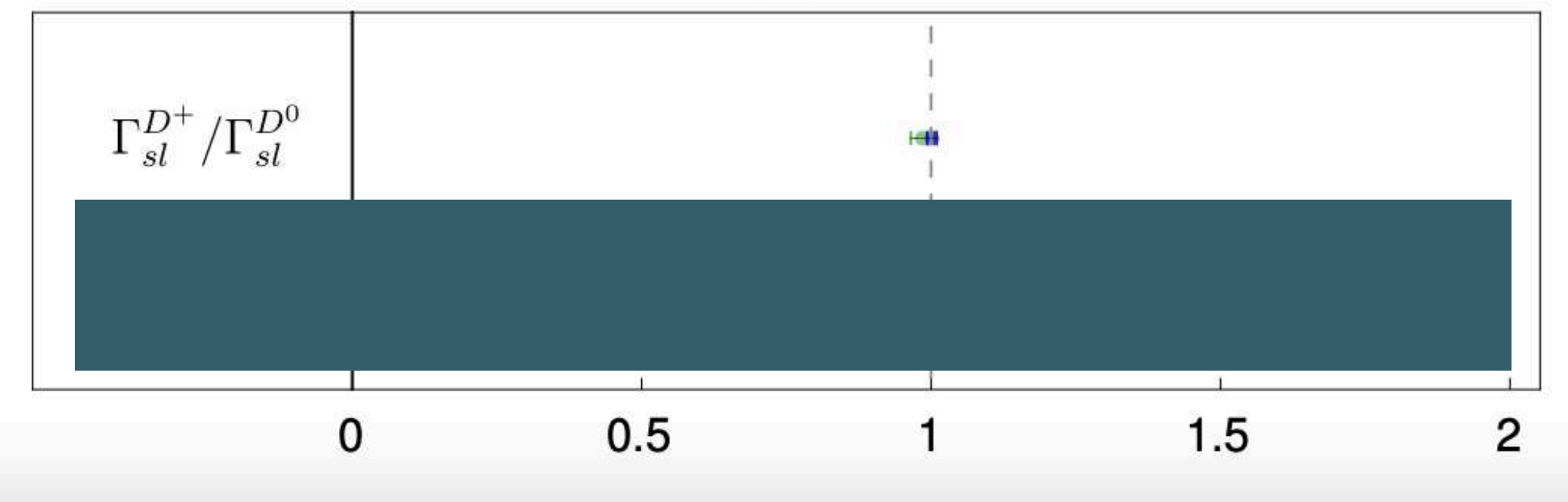
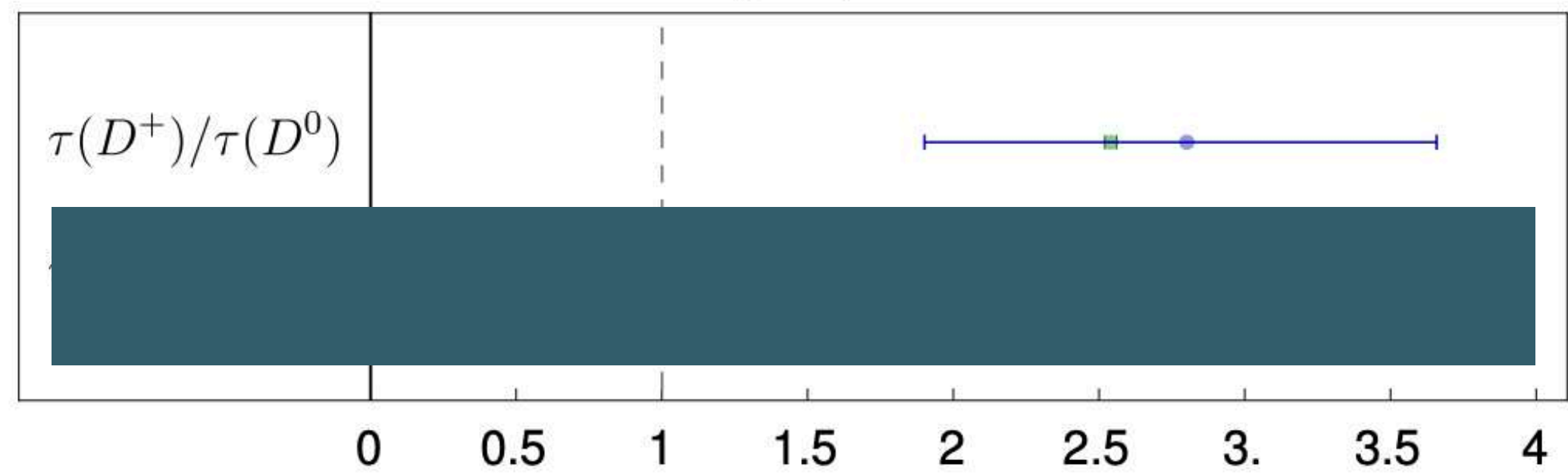
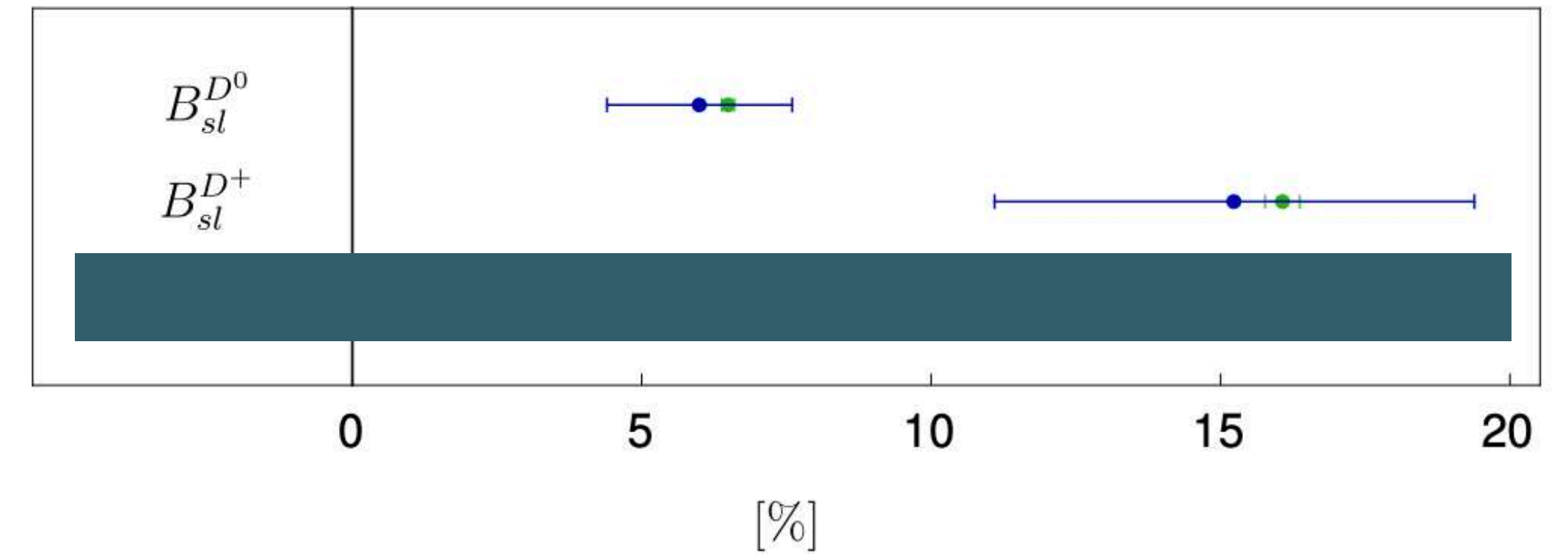
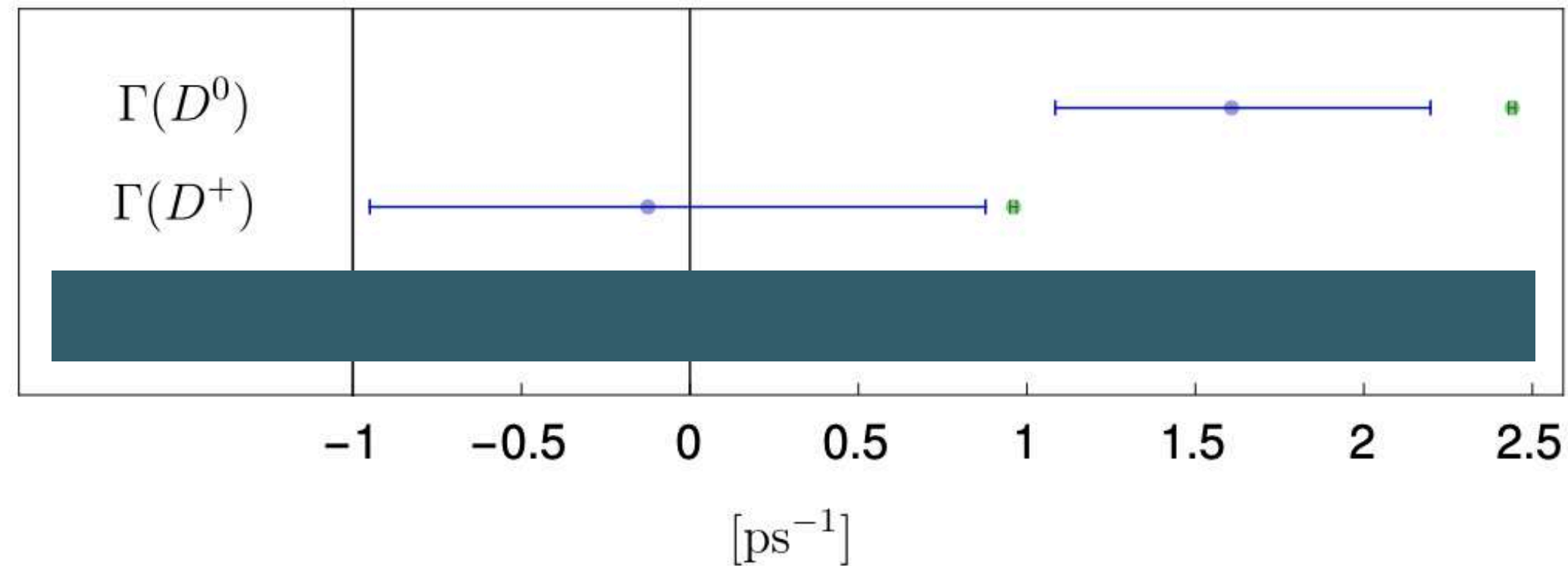
e-Print: 2109.13219 [hep-ph]

$$\Gamma(D^+) = 6.15 \Gamma_0 \left[1 + 0.48 - 0.13 \frac{\mu_\pi^2(D)}{0.48 \text{ GeV}^2} + 0.01 \frac{\mu_G^2(D)}{0.34 \text{ GeV}^2} + 0.31 \frac{\rho_D^3(D)}{0.082 \text{ GeV}^3} \right.$$

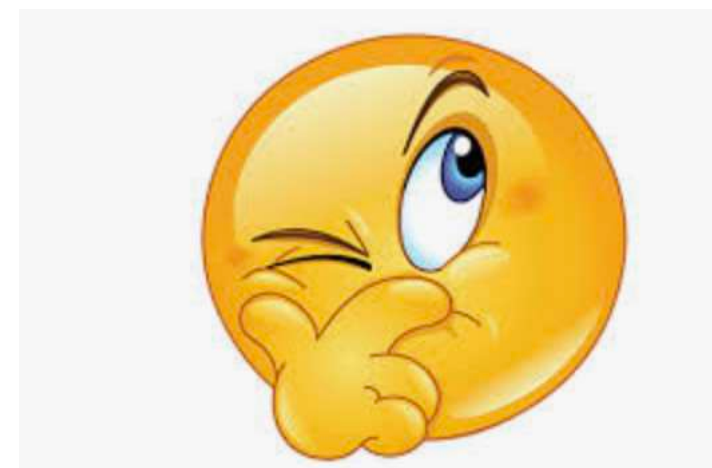


$$\begin{aligned} & - \underbrace{2.66}_{\text{dim-6, VIA}} - 0.055 \frac{\delta \tilde{B}_1^q}{0.02} + 0.002 \frac{\delta \tilde{B}_2^q}{0.02} - 0.546 \frac{\tilde{\epsilon}_1^q}{-0.04} + 0.009 \frac{\tilde{\epsilon}_2^q}{-0.04} + \underbrace{1.10}_{\text{dim-7, VIA}} \\ & - 0.0000 r_1^{qq} - 0.0000 r_2^{qq} + 0.0011 r_3^{qq} + 0.0008 r_4^{qq} \\ & - 0.0109 r_1^{sq} - 0.0080 r_2^{sq} - 0.0000 r_3^{sq} + 0.0001 r_4^{sq} \end{aligned} \Bigg],$$

Huge effects due to Pauli interference



- Values of $\mu_\pi^2, \mu_G^2, \rho_D^3$ almost unknown
- NNLO-QCD corrections to free quark decay in progress
Fael, Steinhauser,...
- NNLO-QCD corrections to spectator effects in progress
Nierste, Steinhauser,...
- Check of HQET sum rule results with lattice
Black, Witzel, ... RBC-UK
- First non-perturbative determination of dimension 7



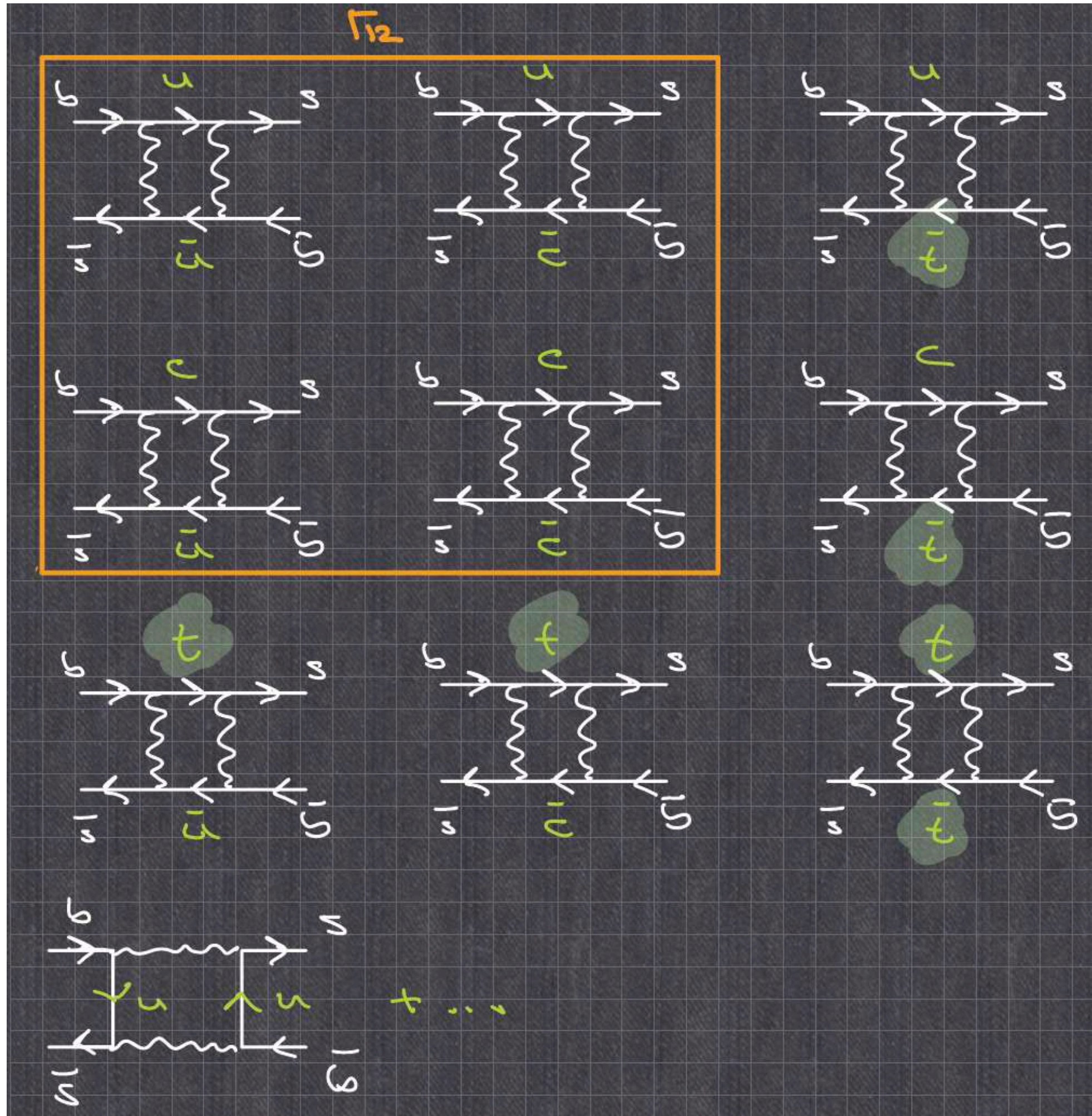
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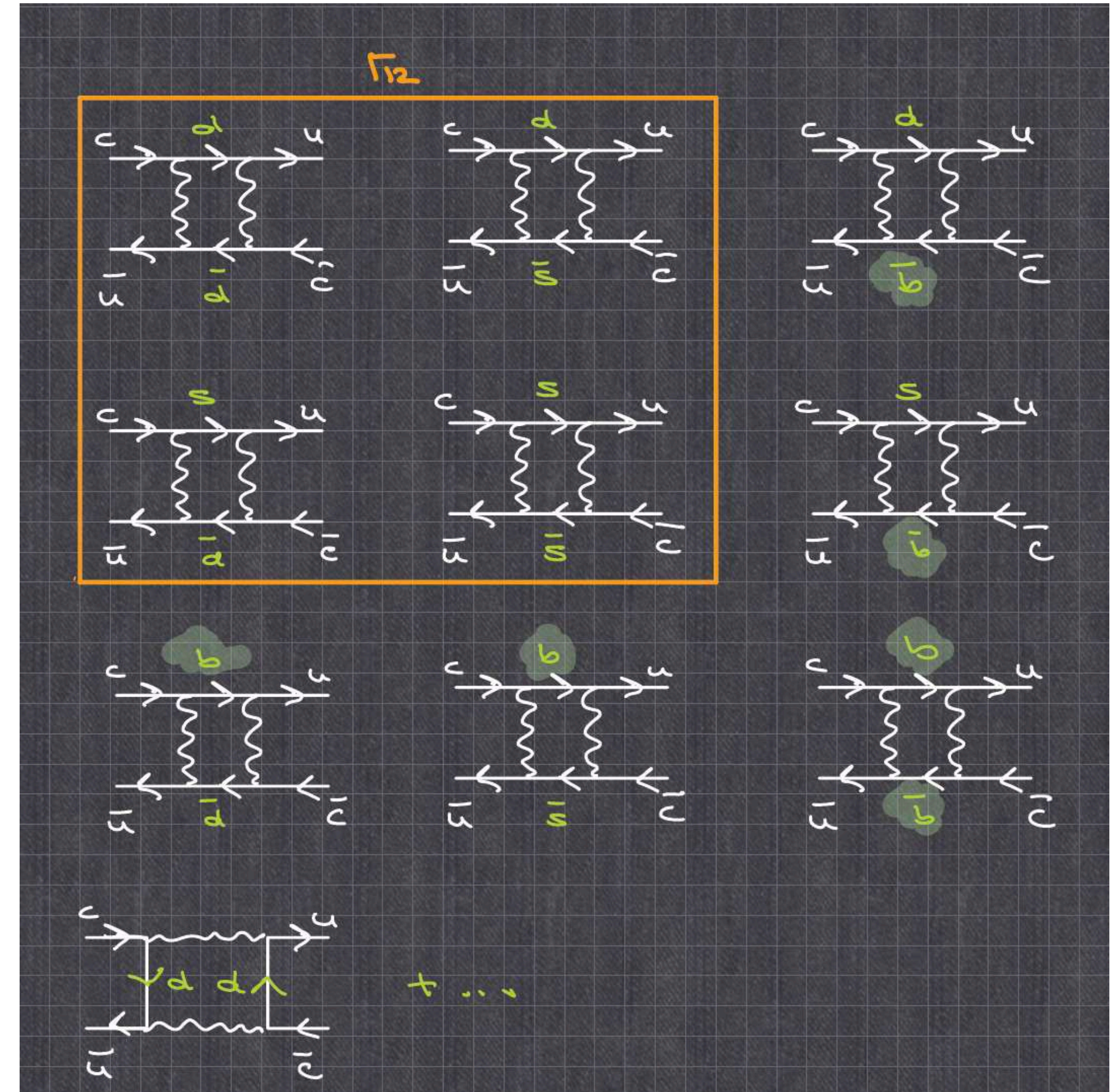
e-Print: 2109.13219 [hep-ph]

C: Crazy Cancellations

B-mixing



D-mixing



C: Crazy Cancellations

B-mixing



$$\begin{aligned}
 M_{12} &= \lambda_u^2 F(u,u) + \lambda_u \lambda_c F(u,c) + \lambda_u \lambda_t F(u,t) \\
 &+ \lambda_c \lambda_u F(c,u) + \lambda_c^2 F(c,c) + \lambda_c \lambda_t F(c,t) \\
 &+ \lambda_t \lambda_u F(t,u) + \lambda_t \lambda_c F(t,c) + \lambda_t^2 F(t,t) \\
 \lambda_u + \lambda_c + \lambda_t &= 0 \\
 &\downarrow \\
 &= \lambda_u^2 [F(c,c) - 2F(u,c) + F(u,u)] \\
 &+ 2\lambda_u \lambda_t [F(c,c) - F(u,c) + F(u,t) - F(c,t)] \\
 &+ \lambda_t^2 [F(c,c) - 2F(c,t) + F(t,t)]
 \end{aligned}$$

	Bd	Bs	
λ_u	$\lambda^{3.8}$	$\lambda^{4.8}$	$w_u^2/w_{\Sigma^2} \approx 0$
λ_c	λ^3	λ^2	$w_c^2/w_{\Sigma^2} \approx 2.5 \cdot 10^{-4}$
λ_t	λ^3	λ^2	$w_t^2/w_{\Sigma^2} \approx 4.5$

CKM dominant \equiv GIM dominant

CKM suppressed \equiv GIM suppressed

D-mixing



$$\begin{aligned}
 M_{12} &= \lambda_d^2 F(d,d) + \lambda_d \lambda_s F(d,s) + \lambda_d \lambda_b F(d,b) \\
 &+ \lambda_s \lambda_d F(s,d) + \lambda_s^2 F(s,s) + \lambda_s \lambda_b F(s,b) \\
 &+ \lambda_b \lambda_d F(b,d) + \lambda_b \lambda_s F(b,s) + \lambda_b^2 F(b,b) \\
 \lambda_d + \lambda_s + \lambda_b &= 0 \\
 &\downarrow \\
 &= \lambda_d^2 [F(d,d) - 2F(d,s) + F(s,s)] \\
 &+ 2\lambda_s \lambda_b [F(s,s) - F(d,s) + F(d,b) - F(s,b)] \\
 &+ \lambda_b^2 [F(s,s) - 2F(s,b) + F(b,b)]
 \end{aligned}$$

	D	
λ_d	λ^1	$w_d^2/w_{\Sigma^2} \approx 0$
λ_s	λ^1	$w_s^2/w_{\Sigma^2} \approx 1.3 \cdot 10^{-6}$
λ_b	$\lambda^{5.8}$	$w_b^2/w_{\Sigma^2} \approx 2.8 \cdot 10^{-3}$

CKM suppressed \equiv GIM dominant

CKM dominant \equiv GIM suppressed

The HQE is successful in the B system and for D meson lifetimes

=> apply it for D-mixing

$$y_D^{\text{HQE}} \approx \lambda_s^2 (\Gamma_{12}^{ss} - 2\Gamma_{12}^{sd} + \Gamma_{12}^{dd}) \approx 10^{-5} y_D^{\text{Exp.}}$$

How can this be?

Look only at a single diagram:

$$y_D^{\text{HQE}} \neq \lambda_s^2 \Gamma_{12}^{ss} \tau_D = 3.7 \cdot 10^{-2} \approx 5.6 y_D^{\text{Exp.}}$$

pert. calculation: **Bobrowski et al 1002.4794**

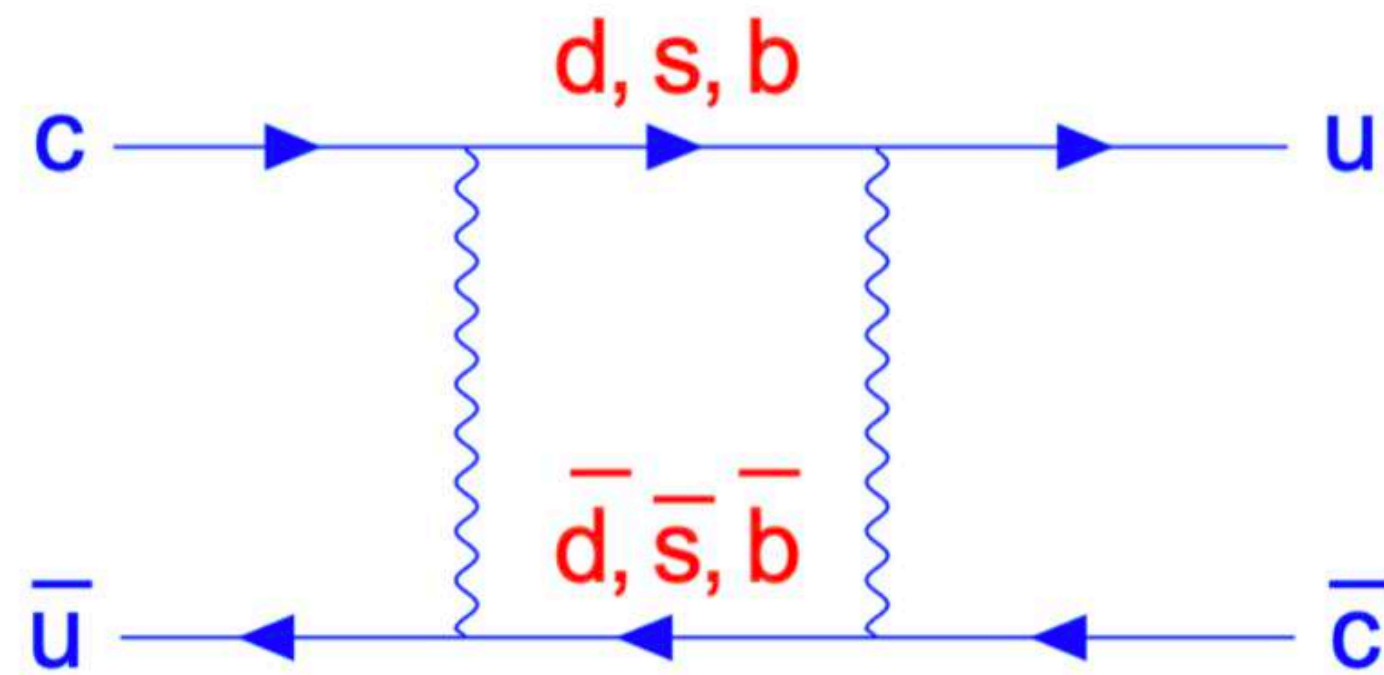
lattice input: **ETM 1403.7302; 1505.06639; FNAL/MILC 1706.04622**

HQET sum rules: **Kirk, AL, Rauh 1711.02100**

The problem seems to originate in the extreme GIM cancellations

C: Crazy Cancellations

GIM cancellation vs CKM hierarchy: $|\lambda_b| \ll |\lambda_s|$, but complex!!!



CPV

survives in
SU(3)F limit!

dominant for
B mixing

$$\Gamma_{12}^D = -\lambda_s^2 (\Gamma_{ss}^D - 2\Gamma_{sd}^D + \Gamma_{dd}^D) + 2\lambda_s\lambda_b (\Gamma_{sd}^D - \Gamma_{dd}^D) - \lambda_b^2 \Gamma_{dd}^D,$$

$$M_{12}^D = \lambda_s^2 [M_{ss}^D - 2M_{sd}^D + M_{dd}^D] + 2\lambda_s\lambda_b [M_{bs}^D - M_{bd}^D - M_{sd}^D + M_{dd}^D] + \lambda_b^2 [M_{bb}^D - 2M_{bd}^D + M_{dd}^D].$$

1. Duality violations - break down of HQE

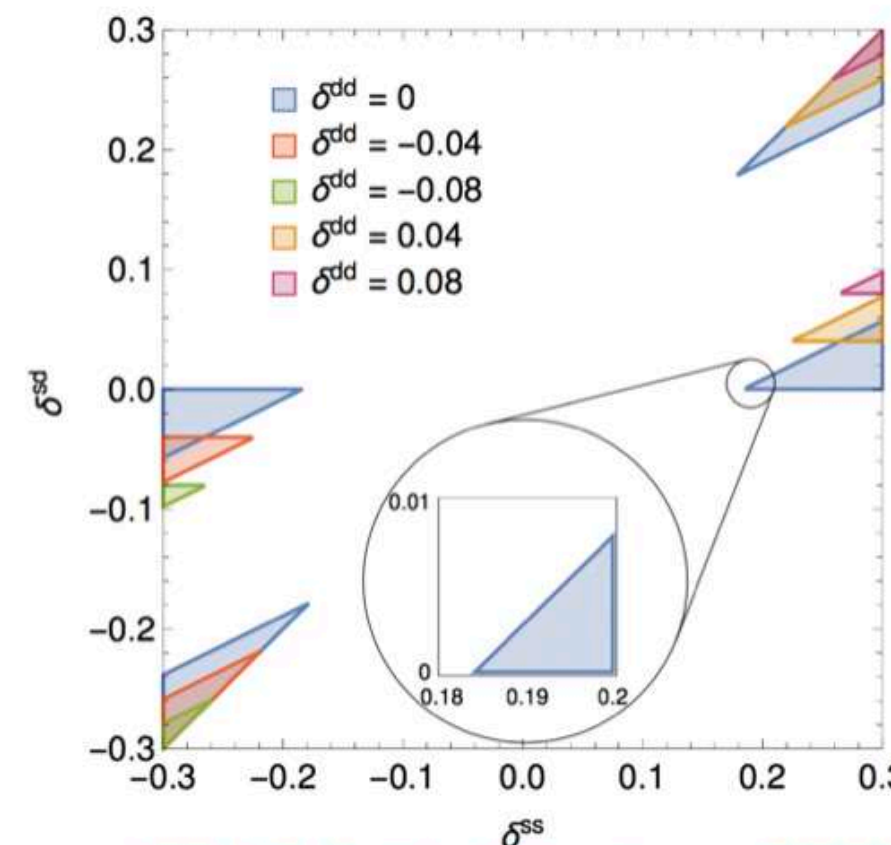
$$\Gamma_{12}^{ss} \rightarrow \Gamma_{12}^{ss}(1 + \delta^{ss}),$$

$$\Gamma_{12}^{sd} \rightarrow \Gamma_{12}^{sd}(1 + \delta^{sd}),$$

$$\Gamma_{12}^{dd} \rightarrow \Gamma_{12}^{dd}(1 + \delta^{dd}),$$

20% of duality violation is sufficient to explain experiment

Jubb, Kirk, AL, Tetlalmatzi-Xolocotzi 2016



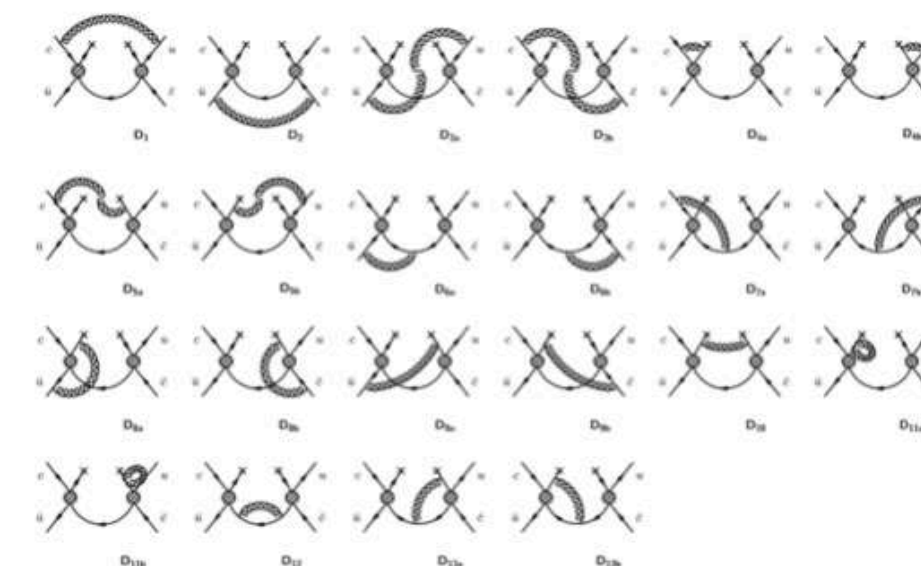
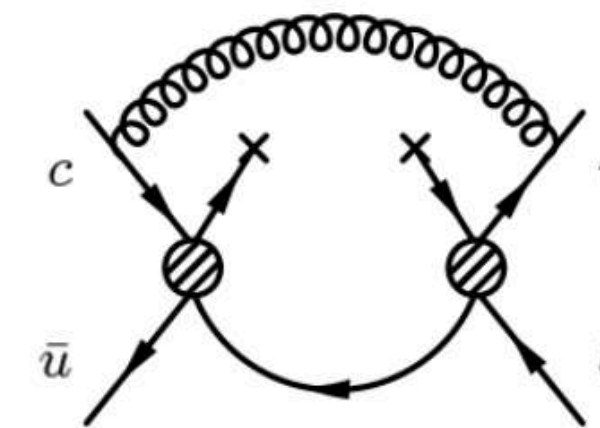
2. Higher dimensions

Georgi 9209291; Ohi, Ricciardi, Simmons 9301212; Bigi, Uraltsev 0005089

Idea: GIM cancellation is lifted by higher orders in the HQE - overcompensating the 1/mc suppression.

Partial calculation of D=9 yields an enhancement - but not to the experimental value

Bobrowski, AL, Rauh 2012



- 1) Vary $\mu^{ss,dd}$ and μ^{ds} independently between 1 GeV and $2 m_c$
 \Rightarrow uncertainty increases and exp. value is covered
- 2) Choose scales somehow phase space inspired as

$$\begin{aligned} \mu^{ss} &= m_c - 2\epsilon \\ \mu^{sd} &= m_c - \epsilon \\ \mu^{dd} &= m_c \end{aligned}$$

\Rightarrow exp. value is covered

3. Renormalisation scale setting:

AL, Piscopo, Vlahos 2020

$$\mu_x^{ss} = \mu_x^{sd} = \mu_x^{dd}$$

Implicitly assumes a precision of 10^{-5} !

4. New Physics is present and we cannot prove it yet:-)

Exclusive and inclusive approaches can cover the experimental regions



No precision determination possible

Exclusive approach

$$\Gamma_{12}^D = \sum_n \rho_n \langle \bar{D}^0 | \mathcal{H}_{eff.}^{\Delta C=1} | n \rangle \langle n | \mathcal{H}_{eff.}^{\Delta C=1} | D^0 \rangle,$$

$$M_{12}^D = \sum_n \langle \bar{D}^0 | \mathcal{H}_{eff.}^{\Delta C=2} | D^0 \rangle + P \sum_n \frac{\langle \bar{D}^0 | \mathcal{H}_{eff.}^{\Delta C=1} | n \rangle \langle n | \mathcal{H}_{eff.}^{\Delta C=1} | D^0 \rangle}{m_D^2 - E_n^2},$$

Cannot be calculated yet

Estimate phase space effects for y : [Falk et al. 0110317](#)

- assume pert. SU(3)_F breaking $y \approx 1\%$
- neglect 3rd family
- neglect SU(3)_F breaking in matrix elements - no QCD calculation

Mass difference from a dispersion relation [Falk et al. 0402204](#) $x \approx y$
 Exp. data [Cheng, Chiang 1005.1106](#) $x \propto \mathcal{O}(0.1\%)$ $y \propto \mathcal{O}(\text{few } 0.1\%)$

U-Spin sum rule [Gronau, Rosner 2012](#)

Factorisation-assisted topological amplitude approach

[Jiang et al. 1705.07335](#) $y \approx 0.2\%$




Direct lattice determination

**Still a very long way!
But not completely crazy
anymore!**

Multiple-channel generalization of Lellouch-Lüscher formula

Maxwell T. Hansen (Washington U., Seattle), Stephen R. Sharpe (Washington U., Seattle) (Apr, 2012)

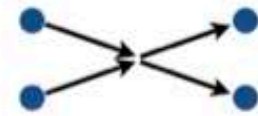

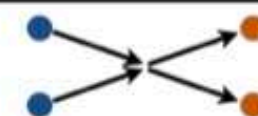
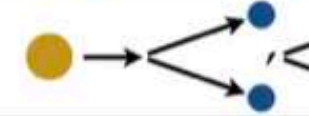
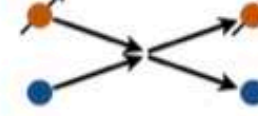

Published in: *Phys.Rev.D* 86 (2012) 016007 • e-Print: 1204.0826 [hep-lat]

 pdf  DOI  cite

#

 259 citation:

Status of multi-hadron matrix elements in LQCD...


physical system	Method to get it from LQCD
$\pi\pi \rightarrow \pi\pi$, $\sqrt{s} < 4M_\pi$ ($\mathbf{P} \neq 0$ in finite-volume frame)*	 Lüscher (1986, 1991) Rummukainen and Gottlieb (1995)*
$K \rightarrow \pi\pi$ (relies on $M_K < 4M_\pi$) ($\mathbf{P} \neq 0$ in finite-volume frame)*	 Lellouch and Lüscher (2001) Kim, Sachrajda and Sharpe (2005)*, Christ, Kim and Yamazaki (2005)*
$\pi\pi \rightarrow K\bar{K}$, $\sqrt{s} < 4M_\pi$ (not possible for physical masses)	 Bernard et al. (2011), Fu (2012), Briceño and Davoudi (2012)
$D \rightarrow \pi\pi, K\bar{K}$ (ignores four-particle states)	 MTH and Sharpe (2012)
$NN \rightarrow NN, N\pi \rightarrow N\pi$ (energies below three-particle production)	 Detmold and Savage (2004) Göckeler et al. (2012) Briceño (2014)
$\gamma^* \rightarrow \pi\pi, \pi\gamma^* \rightarrow \pi\pi,$ $N\gamma^* \rightarrow N\pi$ $B \rightarrow K^*(\rightarrow K\pi)\ell\ell$ (energies below three-particle production)	 Meyer (2011), Bernard et al. (2012), A. Agadjanov et al. (2014), Briceño, MTH and Walker-Loud (2014) Briceño and MTH (2015)

slide by Max Hansen

You Retweeted

Marco Gersabeck (he/him) @MarcoGersabeck · Mar 17

Glad to witness the first seminar of my PhD student Holly McGrath @UoMparticle, which she gave at the #CharmingClues programme at MIAPP. Great job! 🍌 Pity I couldn't join in person.



Charm Decays in Beautiful Places
Semileptonic B_c^+ Decays at LHCb

Adam Davis, Marco Gersabeck, Martha Hilton, Tamaki Holly McGrath

The University of Manchester

MIAPP - Charming Clues for Existence
17th March, 2022

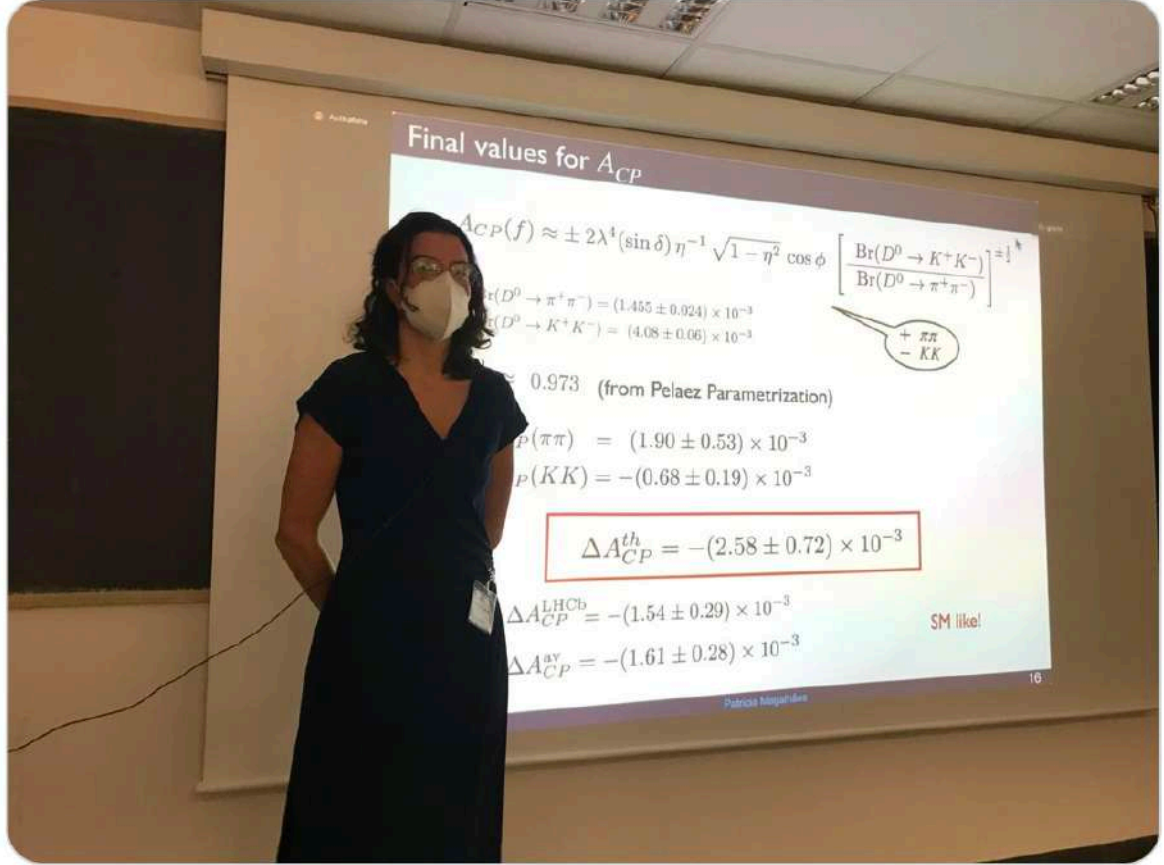
MANCHESTER 1824 LHCb

Tamaki Holly McGrath Semileptonic B_c^+ Decays 17/03/2022 1/21

3 3 30

Alexander Lenz @alexlenz42 · Mar 16

Is CP violation in decays of D mesons a sign of physics beyond the Standard Model? Patricia says no at #charmingclues



Final values for A_{CP}

$$A_{CP}(f) \approx \pm 2\lambda^4 (\sin \delta) \eta^{-1} \sqrt{1 - \eta^2} \cos \phi \left[\frac{\text{Br}(D^0 \rightarrow K^+ K^-)}{\text{Br}(D^0 \rightarrow \pi^+ \pi^-)} \right]^{1/2}$$

$\text{Br}(D^0 \rightarrow \pi^+ \pi^-) = (1.465 \pm 0.024) \times 10^{-3}$
 $\text{Br}(D^0 \rightarrow K^+ K^-) = (4.08 \pm 0.06) \times 10^{-3}$

0.973 (from Pösel Parameterization)

$\Delta A_{CP}(\pi\pi) = (1.90 \pm 0.53) \times 10^{-3}$
 $\Delta A_{CP}(KK) = -(0.68 \pm 0.19) \times 10^{-3}$

$\Delta A_{CP}^{\text{th}} = -(2.58 \pm 0.72) \times 10^{-3}$

$\Delta A_{CP}^{\text{th CB}} = -(1.54 \pm 0.29) \times 10^{-3}$
 $\Delta A_{CP}^{\text{th NP}} = -(1.61 \pm 0.28) \times 10^{-3}$

SM like!

4 11 123

Alexander Lenz @alexlenz42 · Mar 19

OMG! Friday morning at #charmingclues in Munich - what is going on? White Sausages?



3 2 10

Alexander Lenz @alexlenz42 · Mar 10

Is this the charming clue for existence? @MIAPP #charmingclues munich-iapp.de/charmingclues

Alexander Lenz @alexlenz42 · Mar 15

Serious physics discussions about the fundamental laws of physics all night long @MIAPP #charmingclues



3 2 43

You Retweeted

Alexey Petrov @AlexeyPetrov · Mar 12

I gave a colloquium at TUM as part of the "Charming clues for existence" program. I think I'll use this color scheme from now on, what do you think? #charmingclues



Charming colloquium
The ugly side of charm

Table of Contents:

- Charming CP-violation
- Charming mixing
- Charming New Physics

Alexey A. Petrov
Wayne State University

Charming clues for the existence
TUM, 11 March, 2022

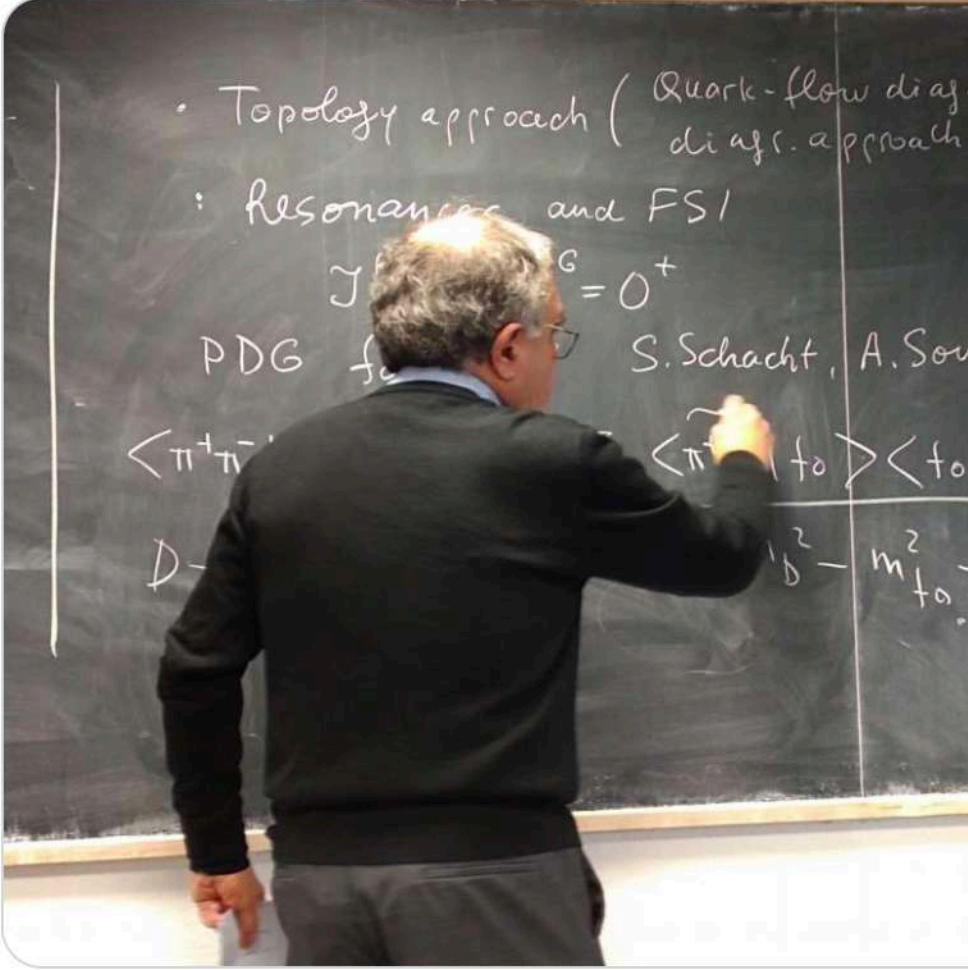
2 10

Alexander Lenz @alexlenz42 · Mar 11

Back to real life meetings with a charming dog, 27 charming colleagues and myself @MIAPP #charmingclues munich-iapp.de/charmingclues



1 5 13



Topology approach (Quark-flow diag. approach)
Resonances and FS1
 $J^P = 0^+$
PDG for S -Schacht, A. Souda
 $\langle \pi^+ \pi^- \rangle$ to $\langle \pi^0 \pi^0 \rangle$
 D - $\langle \pi^+ \pi^- \rangle$ to $\langle \pi^0 \pi^0 \rangle$
 $b^2 - m^2$ to $b^2 - m^2$

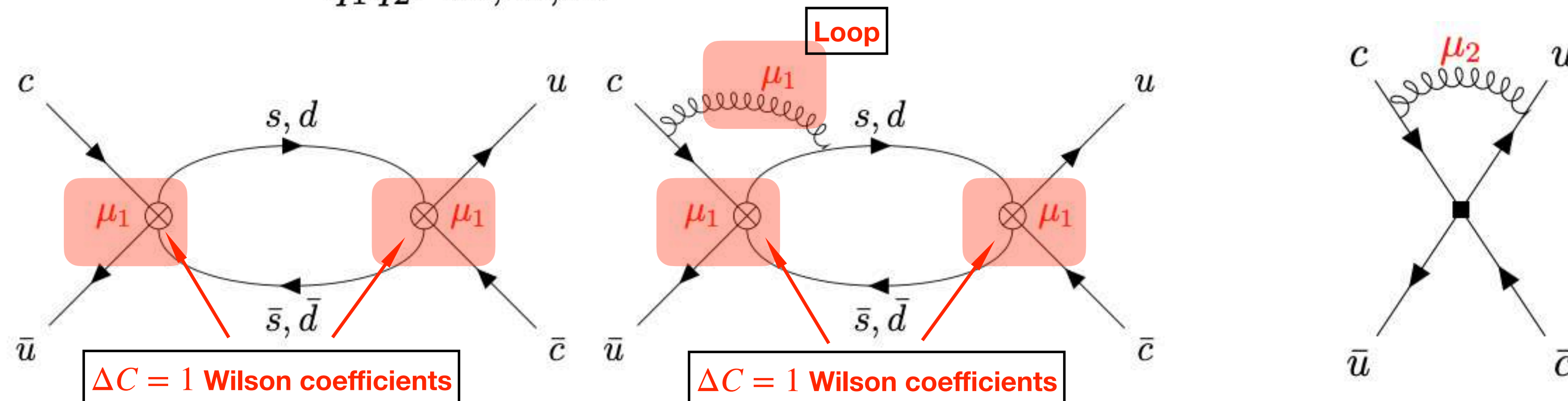
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#Charmingclues

Charm mixing - Theory

Renormalisation scale setting?

$$\Gamma_{12} = \sum_{q_1 q_2 = ss, sd, dd} \Gamma_3^{q_1 q_2} (\mu_1^{q_1 q_2}, \mu_2^{q_1 q_2}) \langle Q \rangle (\mu_2^{q_1 q_2}) \frac{1}{m_c^3} + \dots$$



μ_1 and μ_2 cancel within the ss , sd and dd contributions independently

Is there any requirement to set exactly $\mu_1^{ss} = \mu_1^{sd} = \mu_1^{dd}$ (also during scale variation)?

ss and dd might be related via re-scattering, but sd is physically different from ss !

Charm mixing - Theory

Renormalisation scale setting?

$$\Delta\Gamma_D > 0.028\text{ps}^{-1} \Rightarrow \Omega \equiv \frac{2|\Gamma_{12}|^{\text{SM}}}{0.028\text{ps}^{-1}} \Rightarrow \Omega \approx 1 \text{ means HQE can describe Experiment}$$

Two scenarios:

1. Vary μ^{ss} , μ^{sd} , μ^{dd} independently around m_c between 1 GeV and $2m_c$:

$$\Omega \in [4.6 \cdot 10^{-5}, 1.3]$$

2. Phase space inspired scale choice

$$\mu^{ss} = m_c - 2\epsilon$$

$$\mu^{sd} = m_c - \epsilon$$

$$\mu^{dd} = m_c$$

